

From Molecular Clouds to Massive Stars Star Formation in Numerical Simulations

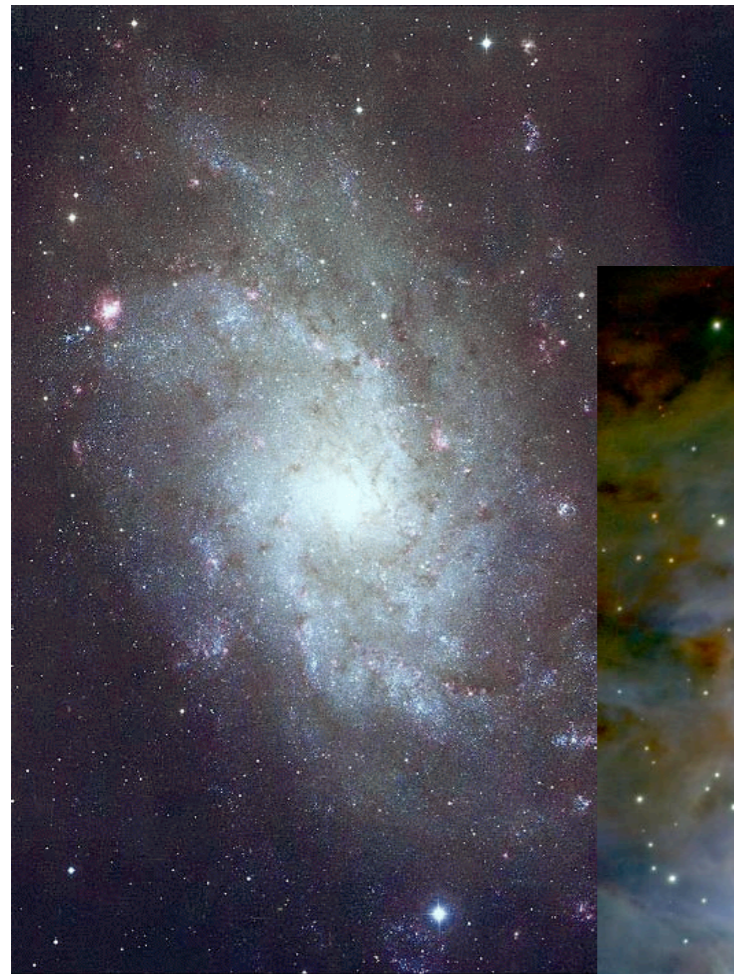
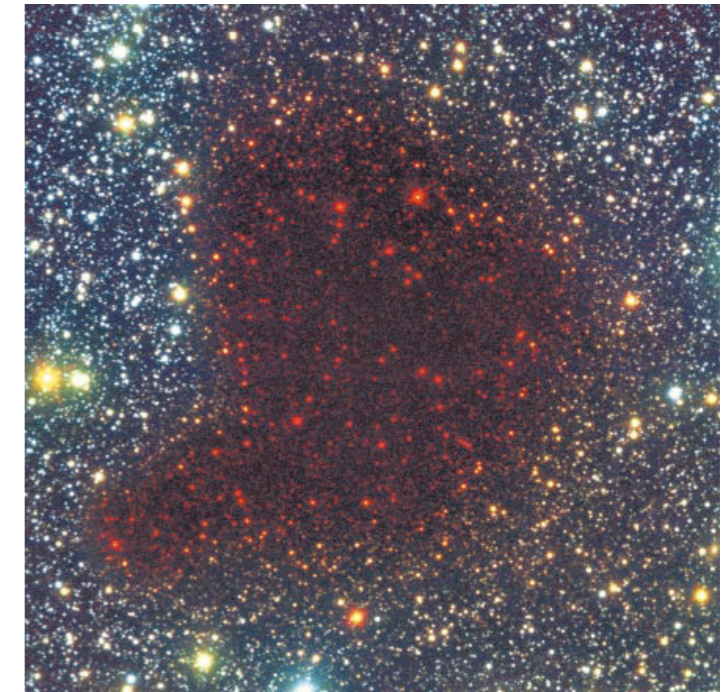
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Collaborators:

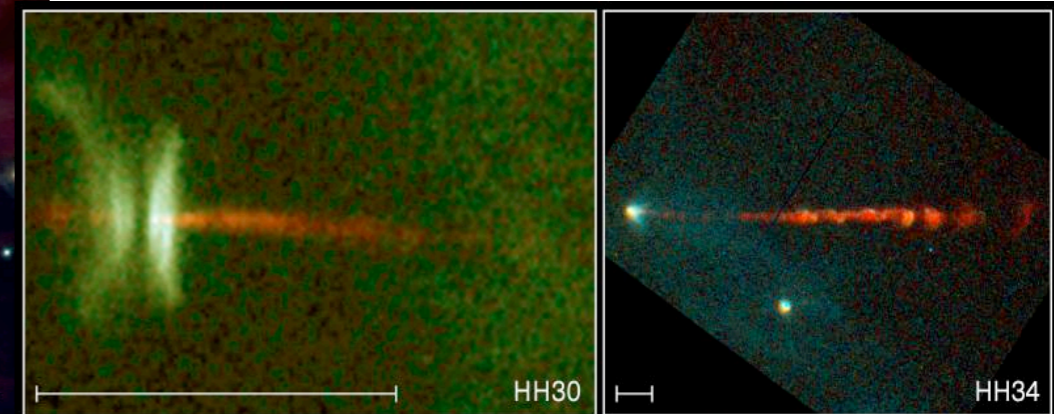
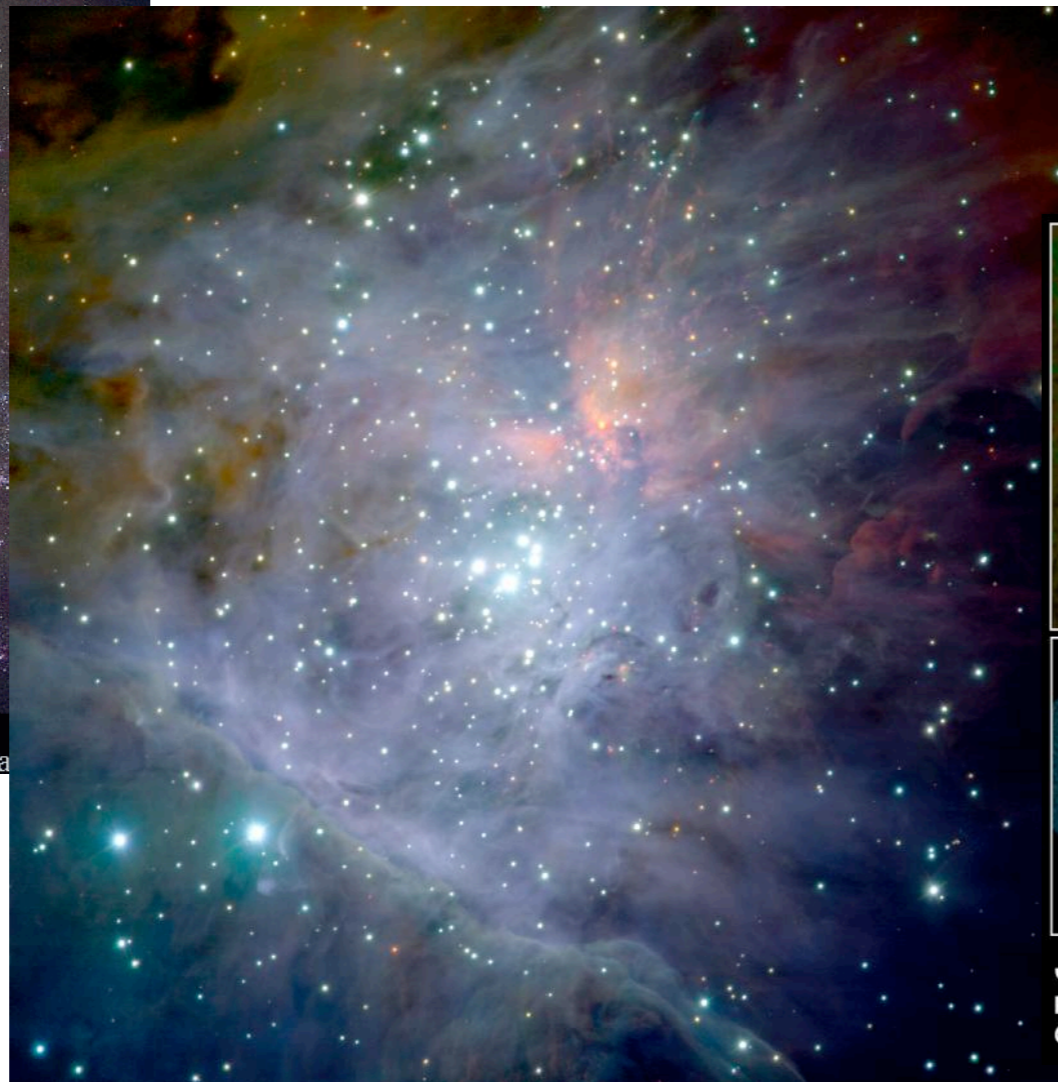
Ralf Klessen, Paul Clark, Christoph Federrath, Philipp Girichidis, Thomas Peters (ITA)
Ralph Pudritz, Dennis Duffin (McMaster), Enrique Vazquez-Semadeni, Roberto Galvan-Madrid (UNAM),
Patrick Hennebelle (ENS), Eric Keto (CfA)

Motivation

Complexity of physical processes (gravity, turbulence, feedback) requires **numerical simulations** to study star formation



M33 © IAC/RGO/Malin
Photo from Isaac Newton Telescope plates by Da



Jets from Young Stars
PRC95-24a · ST Scl OPO · June 6, 1995
C. Burrows (ST Scl), J. Hester (AZ State U.), J. Morse (ST Scl), NASA

HST · WFPC2

Formation of Molecular Clouds

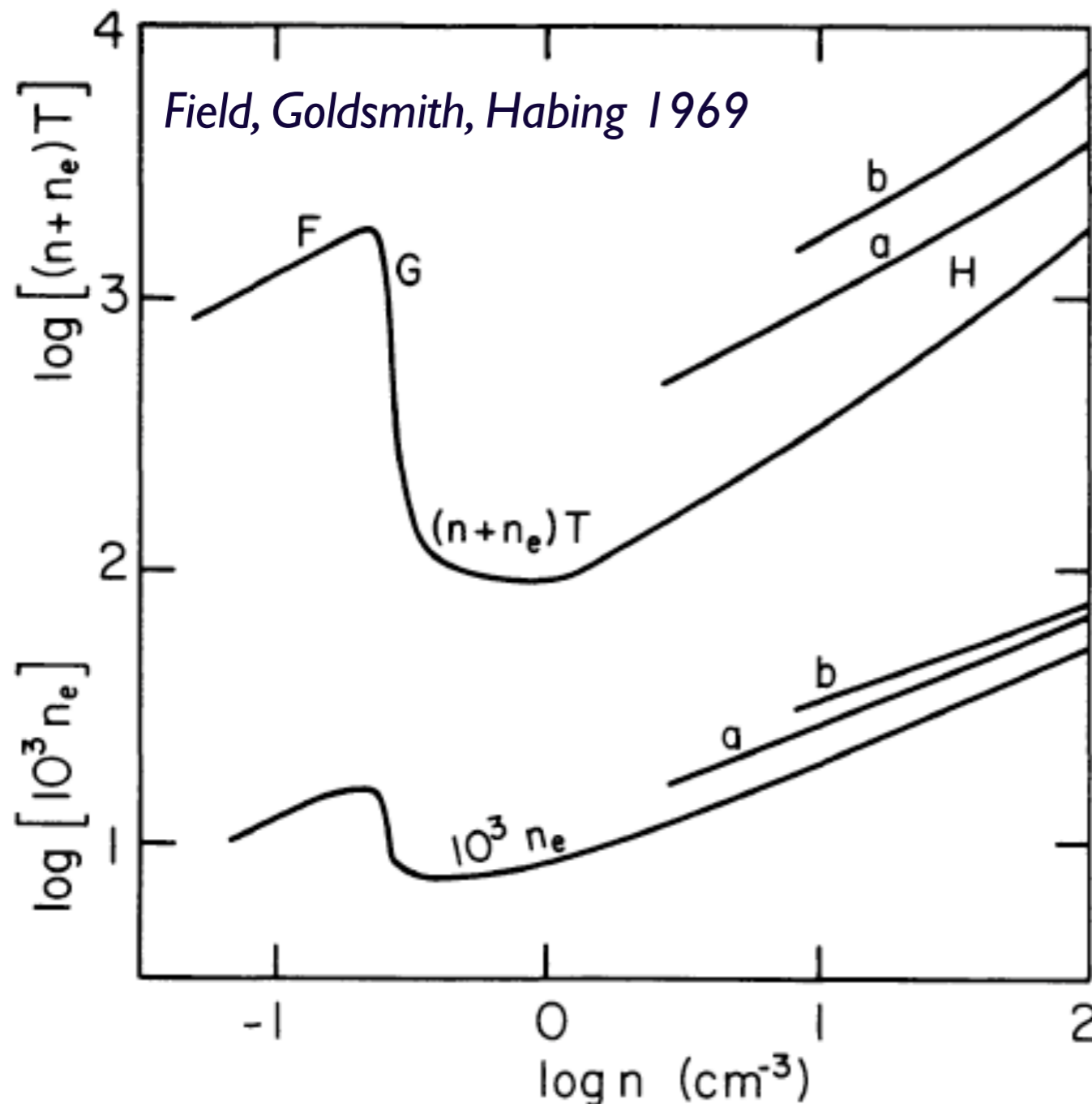
- ISM: warm ionized gas, warm atomic gas, cold molecular gas, dust
- Irregular, inhomogeneous distribution
- highly turbulent
- magnetized
- Present day star formation happens in **Giant Molecular Clouds**
- properties of GMC: density $\sim 10^{2-3} \text{ cm}^{-3}$, size \sim tens pc, mass $\sim 10^{4-6} M_{\text{sol}}$
- composition: 70% hydrogen, 1% dust (mass)
- e.g. nearby Orion nebula (d \sim 400 pc)



M i l k y W a y G a l a x y

Thermal Instability

Formation of dense, cold clouds out of the warm medium through **thermal instability** (Field 1965)?



$$\frac{\partial \ln p}{\partial \ln \rho} < 0 \quad \text{necessary condition for TI}$$

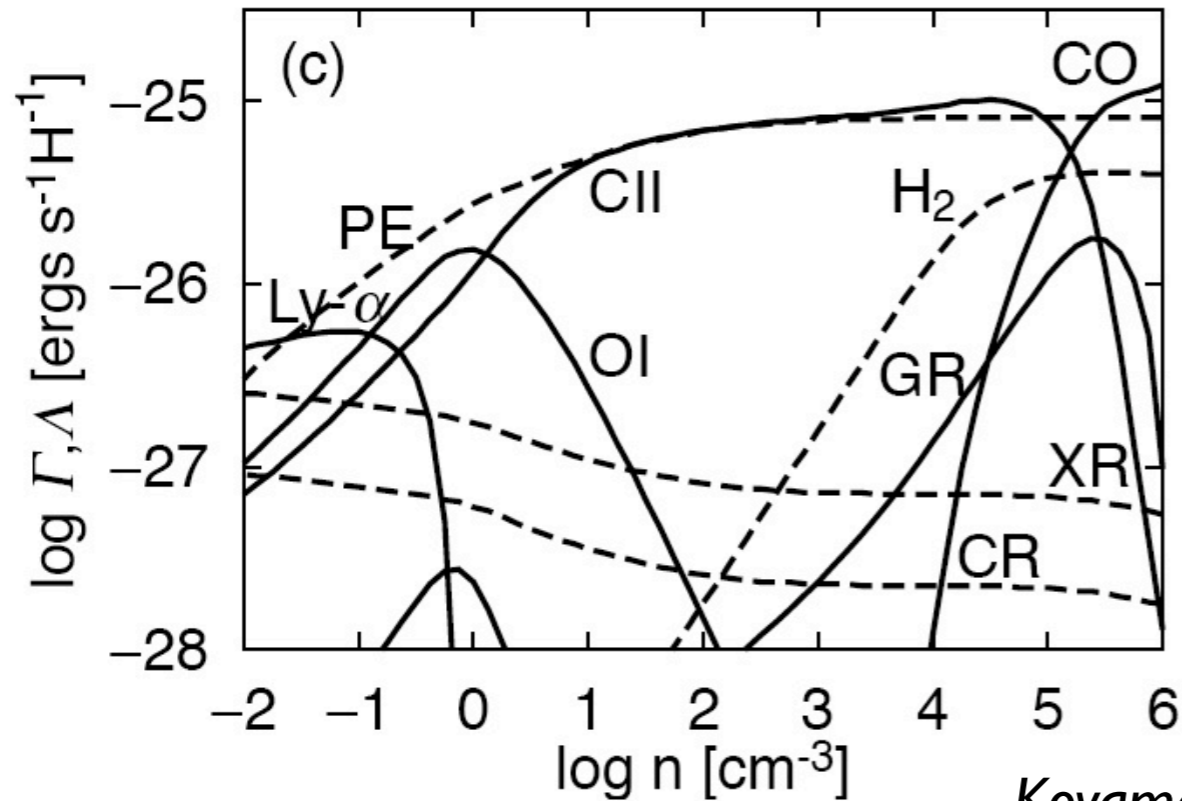
heating (UV, cosmic ray) and **cooling** (atomic and molecular line emission, gas-dust coupling) regulate thermodynamics

Note:

$$M_{\text{Jeans}}(\text{warm gas}) \gg M_{\text{cloud}}$$

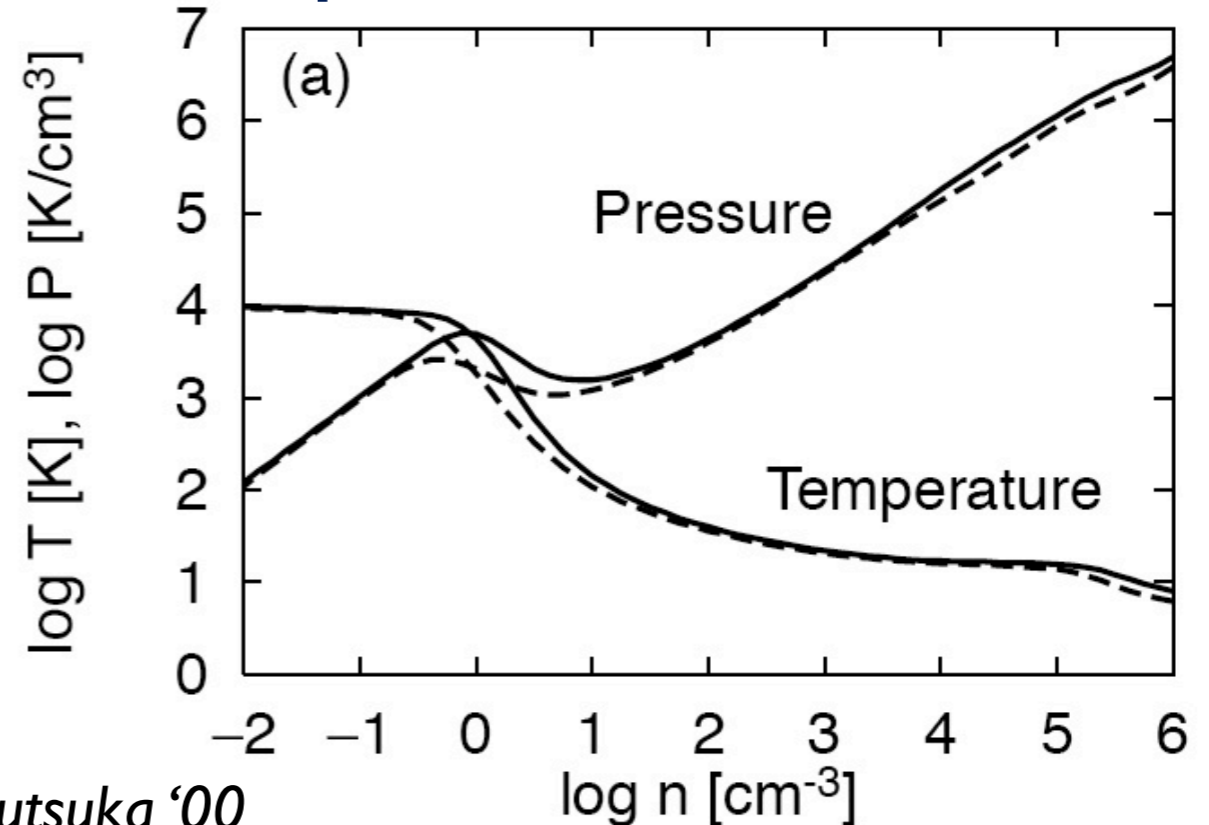
Thermal Instability

main cooling & heating processes



Koyama & Inutsuka '00

equilibrium pressure / temperature



simplification:

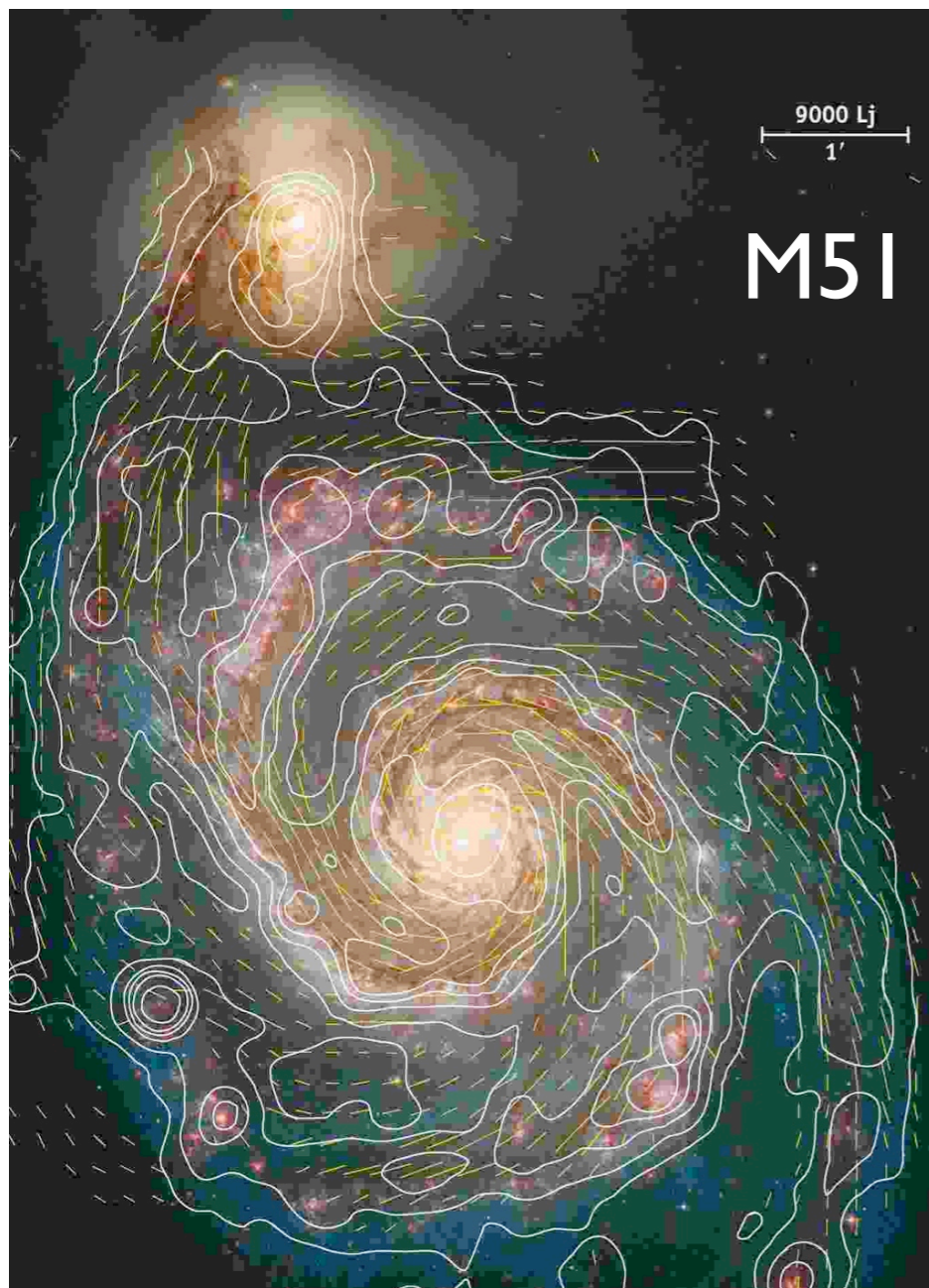
$$\Gamma = 2.0 \times 10^{-26} \text{ ergs}^{-1},$$

$$\frac{\Lambda(T)}{\Gamma} = 10^7 \exp\left(\frac{-1.184 \times 10^5}{T + 1000}\right) + 1.4 \times 10^{-2} \sqrt{T} \exp\left(\frac{-92}{T}\right) \text{ cm}^3$$

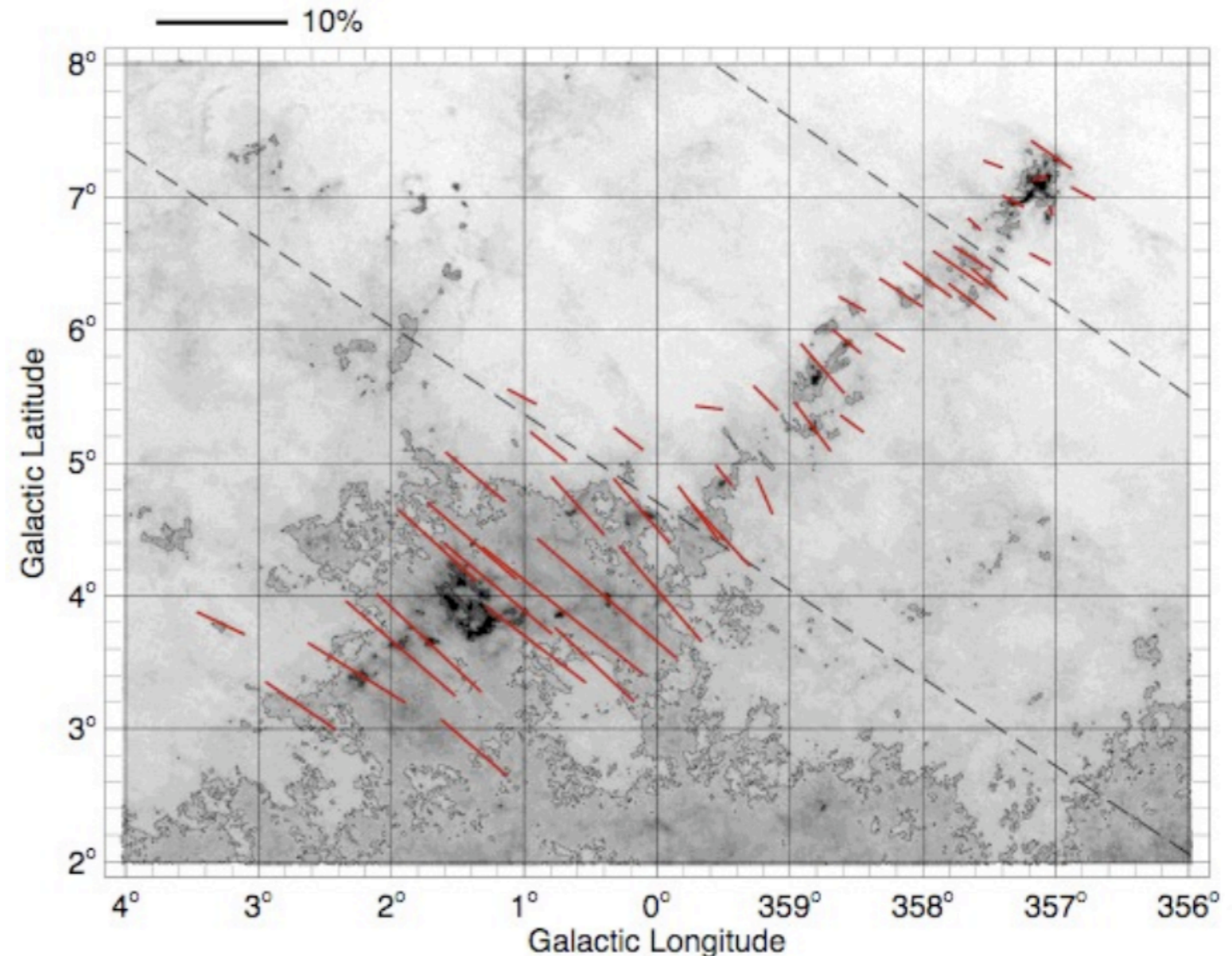
Koyama & Inutsuka '02

Magnetic Fields

The ISM is permeated with magnetic fields



galactic B-fields (e.g. *R.Beck 2001*)
large scale component: $\sim 4\mu\text{G}$
total field strength: $\sim 6\mu\text{G}$



magnetic polarization measurements in the Pipe nebula
F.O.Alves, Franco, Girart 2008

Magnetic Fields

magnetic criticality

mass-to-flux ratio:

$$\mu \equiv \left(\frac{M}{\Phi} \right) = \text{self-gravity} / \text{magnetic support}$$

critical value:

$$\mu_{\text{crit}} = \frac{1}{2\pi \sqrt{G}} \approx 0.16 / \sqrt{G}$$

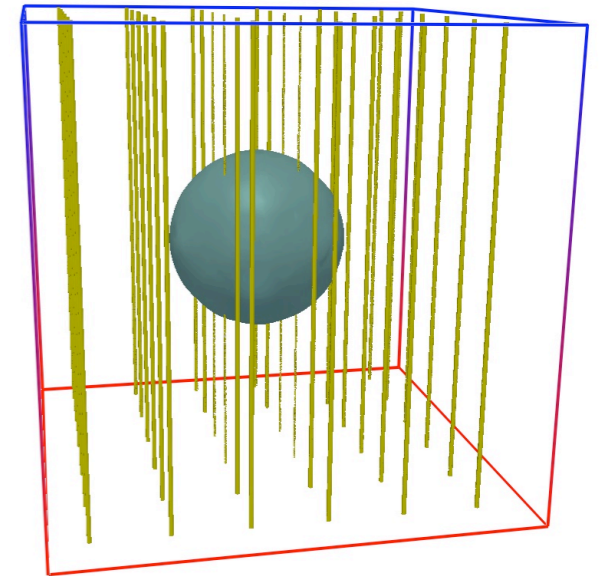
uniform disc

Nakano & Nakamura 1978

$$\mu_{\text{crit}} = 0.13 / \sqrt{G}$$

flattened collapsing structure

Mouschovias & Spitzer 1976

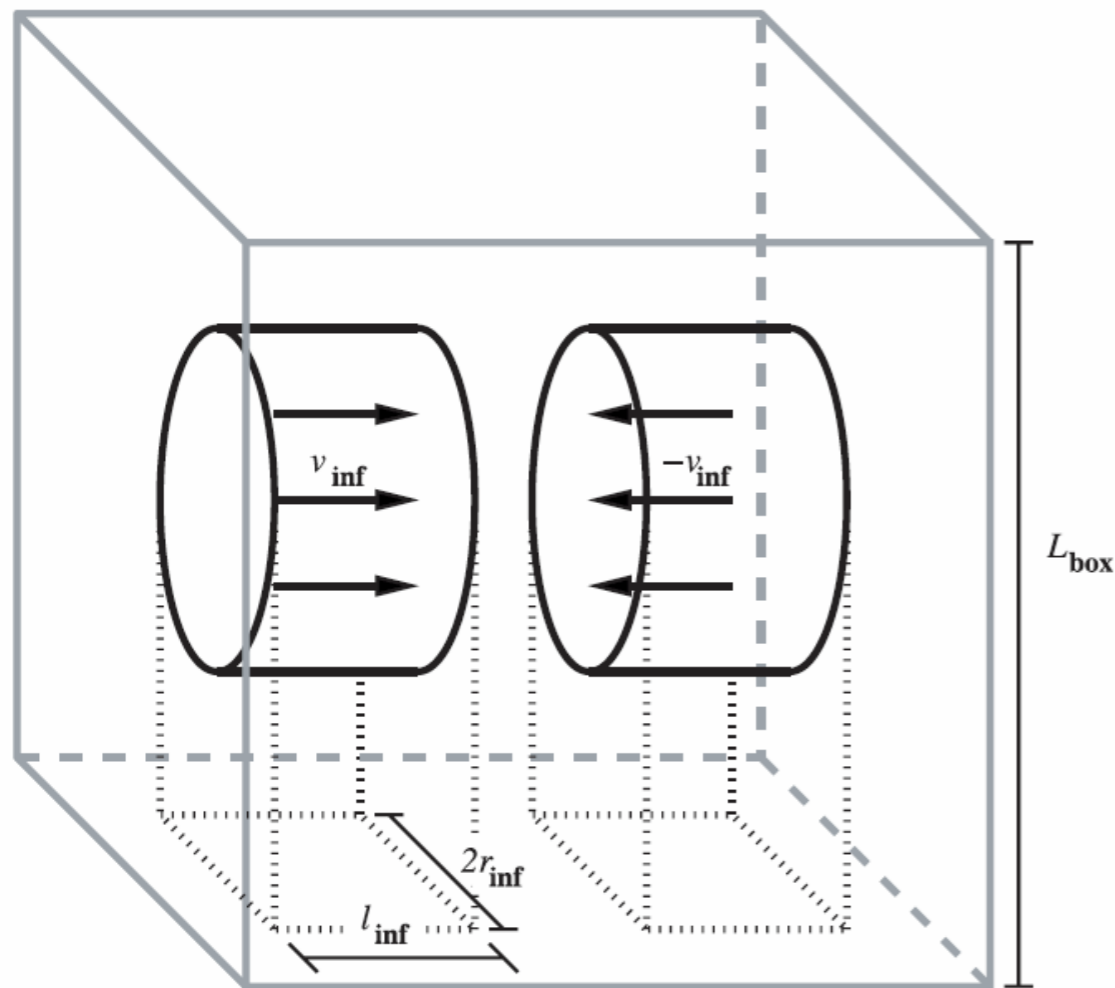


but: Ambipolar diffusion allows sub-critical cores to collapse

Formation of Molecular Clouds

3D simulations with FLASH*

Large scale converging flows



Vazquez-Semadeni et al. 2007

Fiducial model parameter:

- $L_{\text{box}} = 256 \text{ pc}$, $\Delta x_{\text{min}} = 0.03 \text{ pc}$
- $l_{\text{inf}} = 112 \text{ pc}$
- $r_{\text{inf}} = 32 \text{ pc}$
- $v_{\text{inf}} = 6.9 \text{ km/sec} = 1.22 M_a$
- $n = 1 \text{ cm}^{-3}$
- $M_{\text{inf}} = 2.3 \times 10^4 M_{\text{sol}}$
- $T = 5000 \text{ K}$
- $M_J = 10^7 M_{\text{sol}}$
- $B = 1 \mu\text{G}$ aligned with the flow
- $\mu = 3.3$ (super-critical)

*Alliance Center for Astrophysical
Thermonuclear Flashes (ASC), University of
Chicago
Current Version: 3.1

Numerical Method

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\mathbf{v} \rho) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B}) + \nabla p_* = -\rho \mathbf{g}$$

$$\frac{\partial \rho E}{\partial t} + \nabla \cdot (\mathbf{v} (\rho E + p_*) - \mathbf{B} (\mathbf{v} \cdot \mathbf{B})) = \rho \mathbf{g} \cdot \mathbf{v} + \Gamma - \Lambda$$

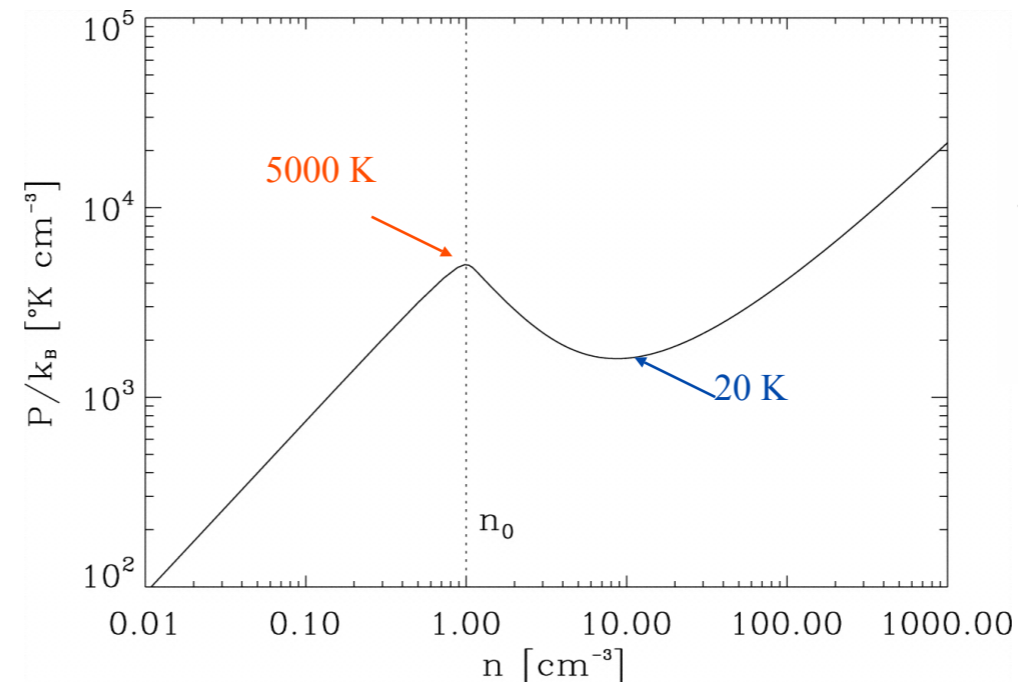
$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v}) = 0$$

$$E = \frac{1}{2} v^2 + \epsilon + \frac{1}{2} \frac{B^2}{\rho},$$

$$p_* = p + \frac{B^2}{2},$$

$$p = (\gamma - 1) \rho \epsilon$$

$$\mathbf{g} = -\nabla \Phi \quad \Delta \Phi = 4\pi G \rho$$



$$\Gamma = 2.0 \times 10^{-26} \text{ ergs}^{-1},$$

$$\frac{\Lambda(T)}{\Gamma} = 10^7 \exp\left(\frac{-1.184 \times 10^5}{T + 1000}\right)$$

$$+ 1.4 \times 10^{-2} \sqrt{T} \exp\left(\frac{-92}{T}\right) \text{ cm}^3$$

Koyama & Inutsuka '02

Ideal MHD + self-gravity + ideal gas + heating & cooling

Formation of Molecular Clouds

the non-magnetic case

edge-on view

face-on view

Formation of Molecular Clouds

the non-magnetic case

0.00 Myr

Boxsize 80.0 pc

edge-on view

0.00 Myr

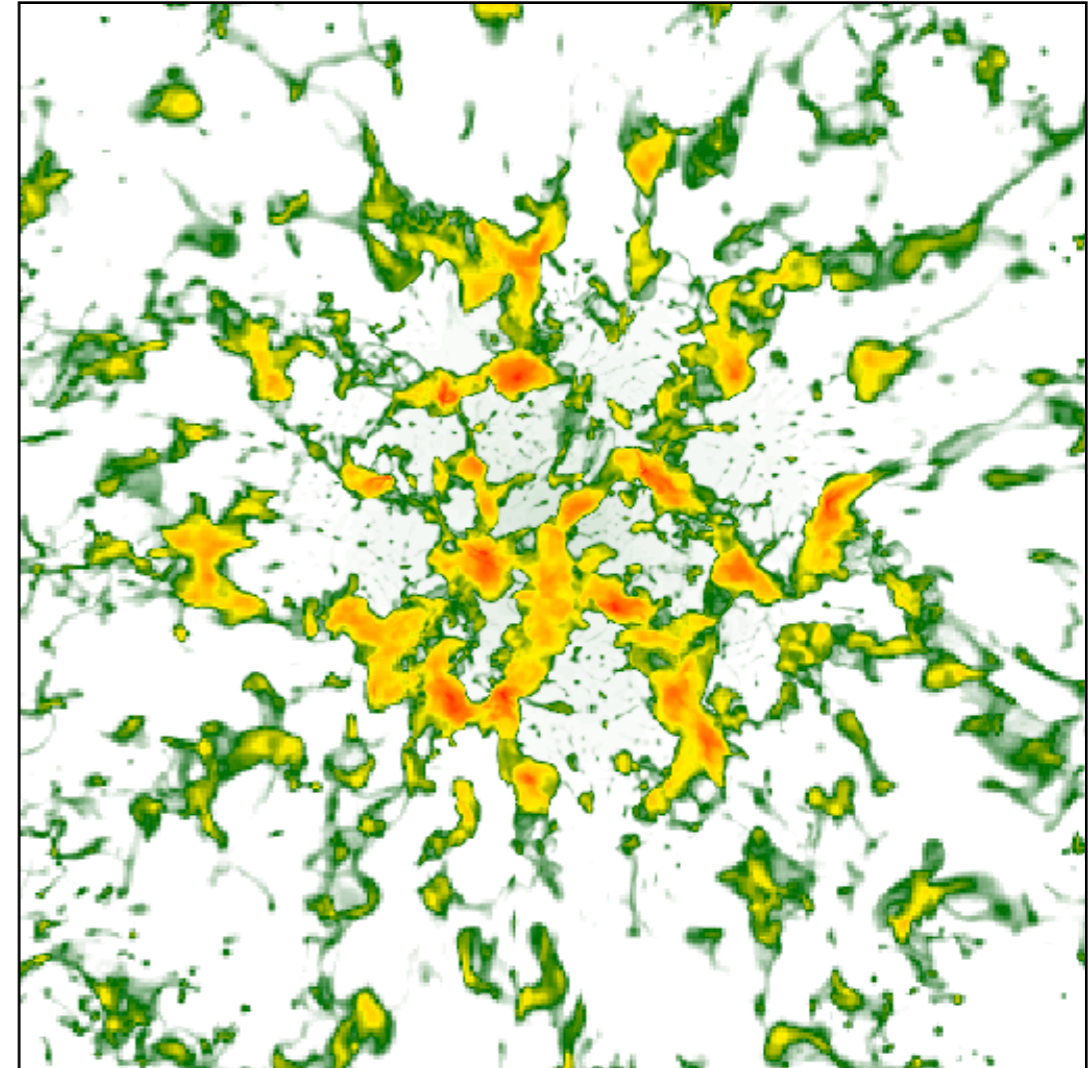
Boxsize 80.0 pc

face-on view

Formation of Molecular Clouds

main properties of MCs:

- highly **patchy** and **clumpy**
- high fraction of **substructure**
- cold dense molecular clumps **coexist** with warm atomic gas
- not a well bounded entity
- **dynamical** evolution (different star formation modes: from low mass to high mass SF?)



Formation of Molecular Clouds

the weakly magnetized ($B_x = 1 \mu\text{G}$) case

0.00 Myr

Boxsize 80.0 pc

edge-on view

0.00 Myr

Boxsize 80.0 pc

face-on view

Formation of Molecular Clouds

with random component: $B_x = 3\mu G + \delta b = 3\mu G$

face-on view

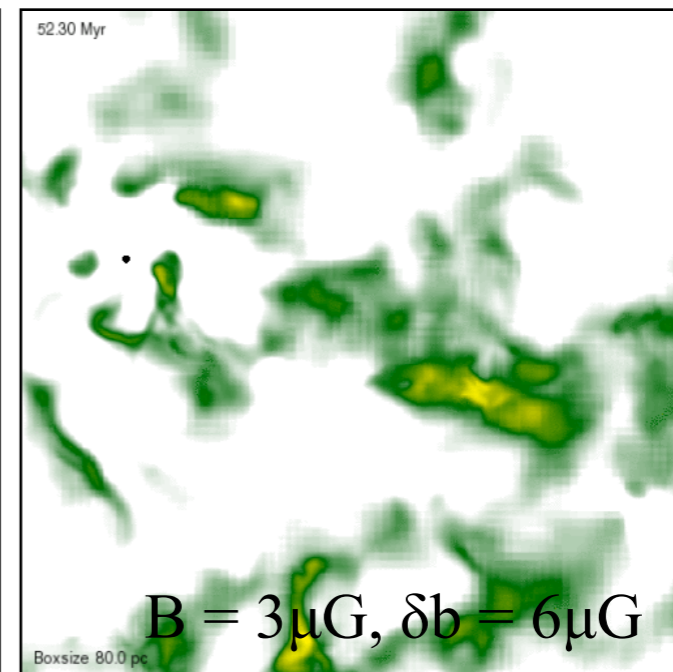
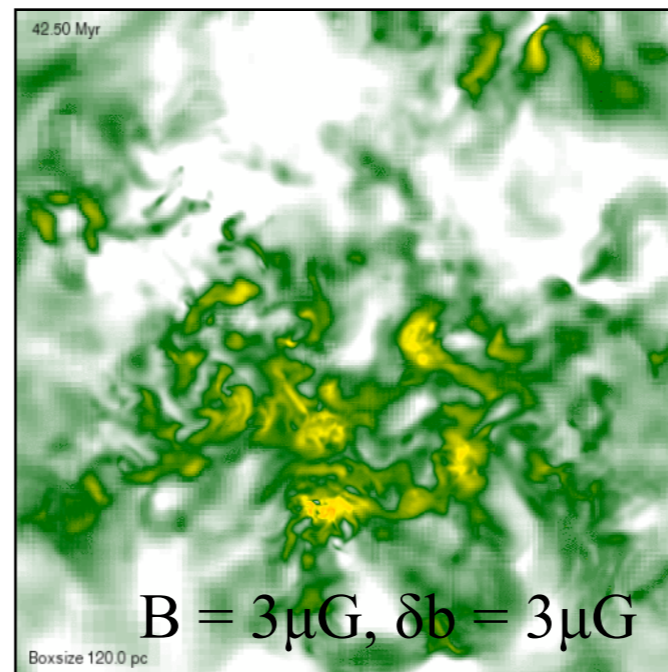
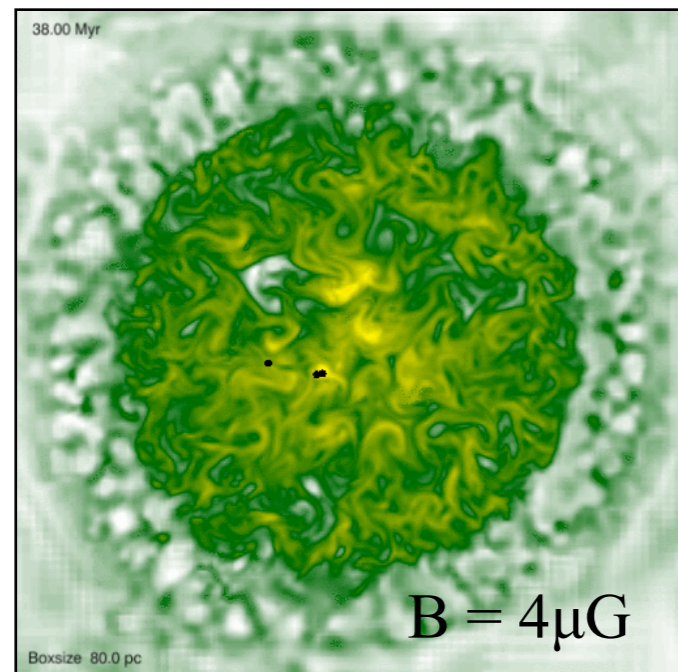
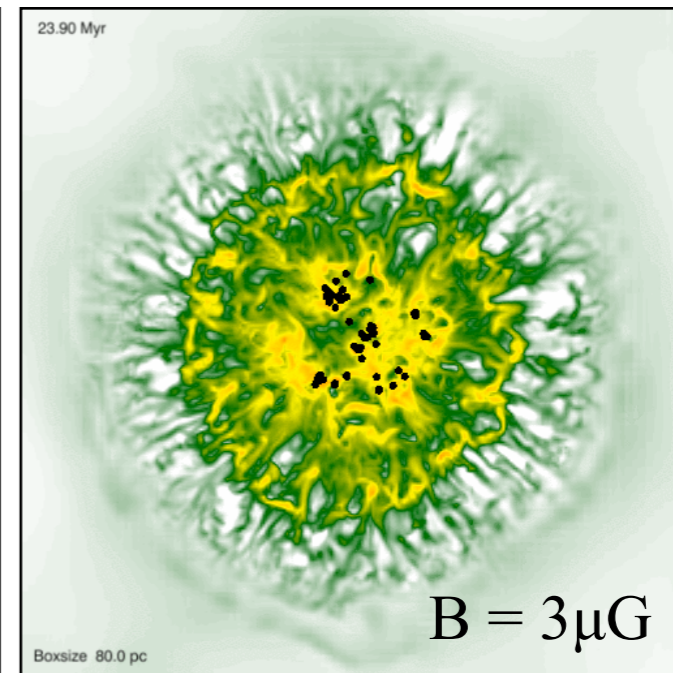
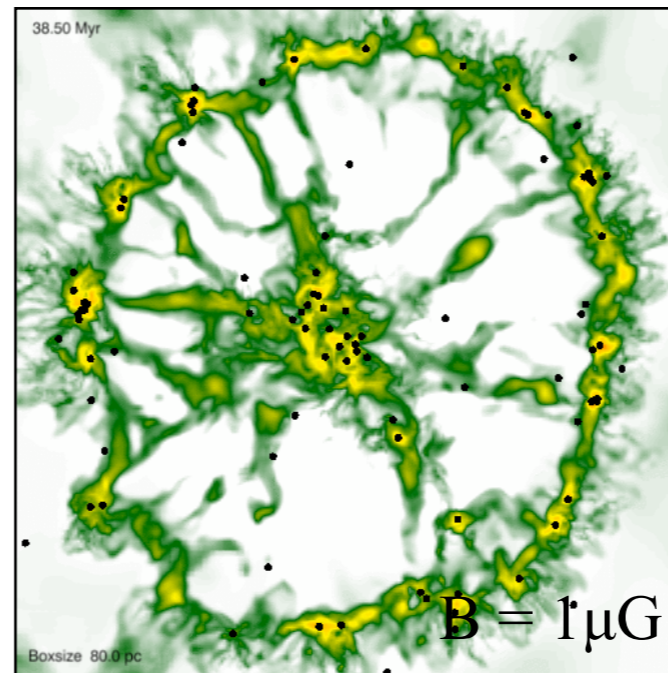
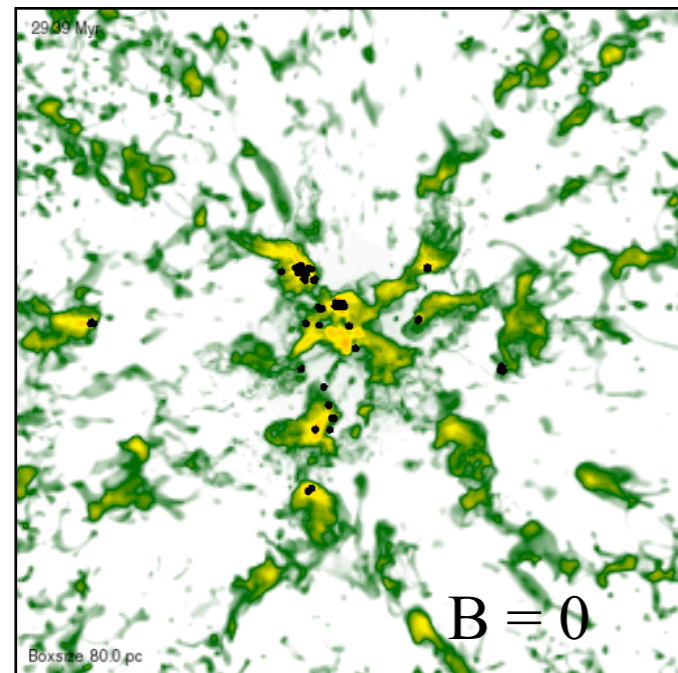
Formation of Molecular Clouds

with random component: $B_x = 3\mu\text{G} + \delta b = 3\mu\text{G}$



face-on view

Formation of Molecular Clouds



Morphology of the molecular cloud and **star formation** efficiency depends on the strength of the magnetic field

Formation of Molecular Clouds

Influence of Ambipolar Diffusion: $B_x = 3\mu\text{G}$ (super-critical)

0.00 Myr

Boxsize 80.0 pc

Ideal MHD

0.00 Myr

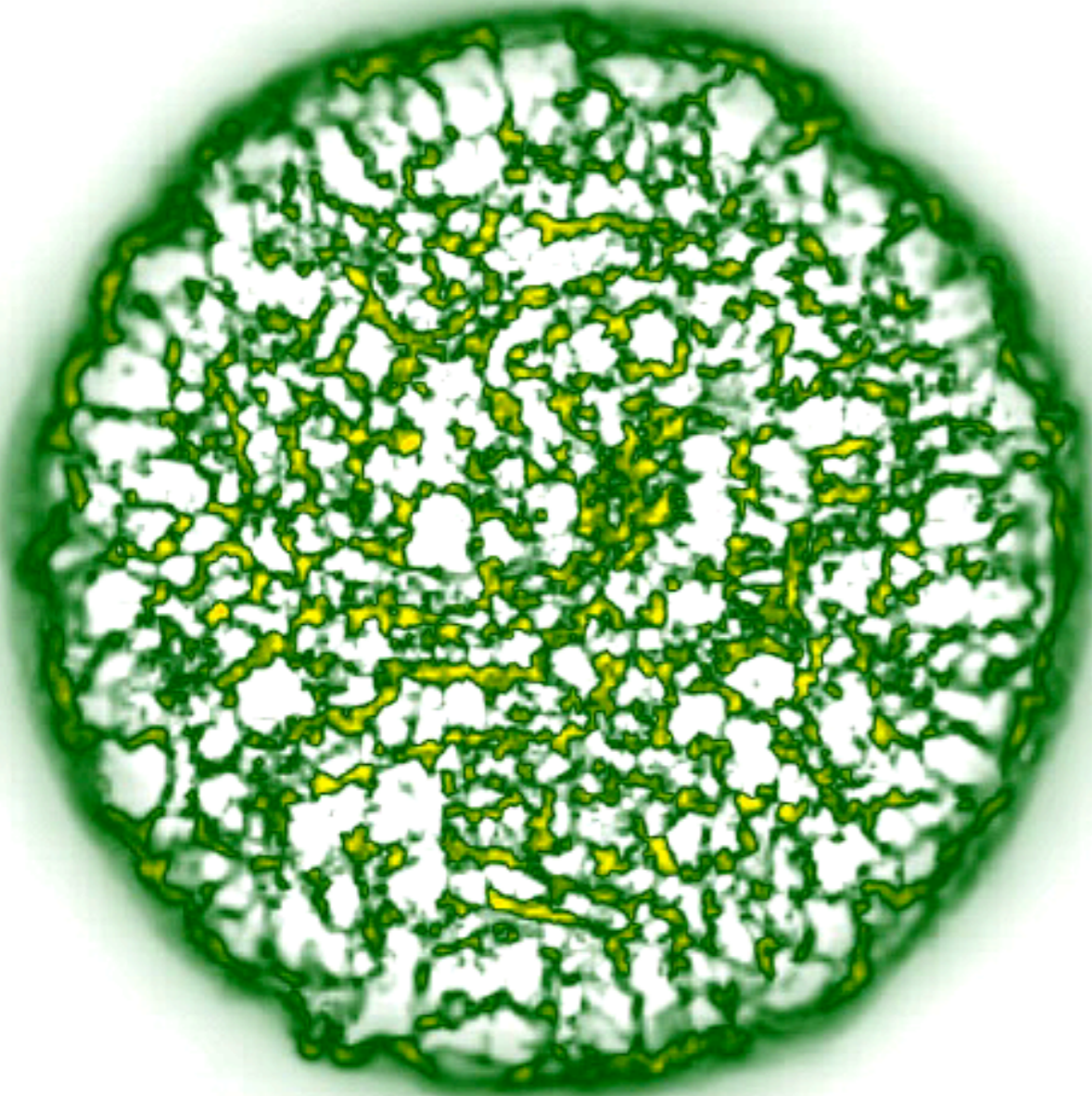
Boxsize 80.0 pc

with AD

Formation of Molecular Clouds

Influence of Ambipolar Diffusion: $B_x = 4\mu\text{G}$ (critical)

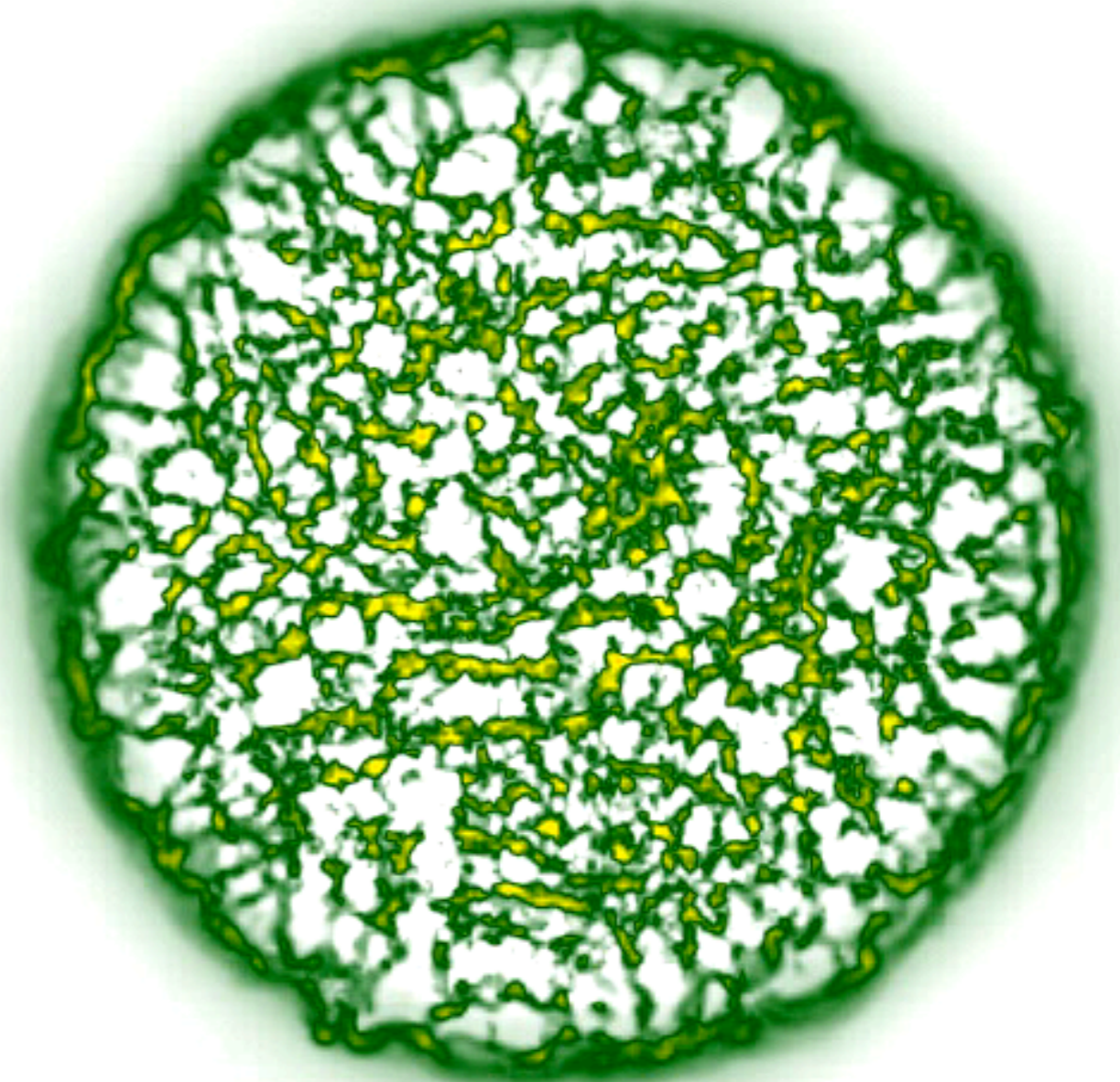
7.00 Myr



Boxsize 80.0 pc

Ideal MHD

6.90 Myr

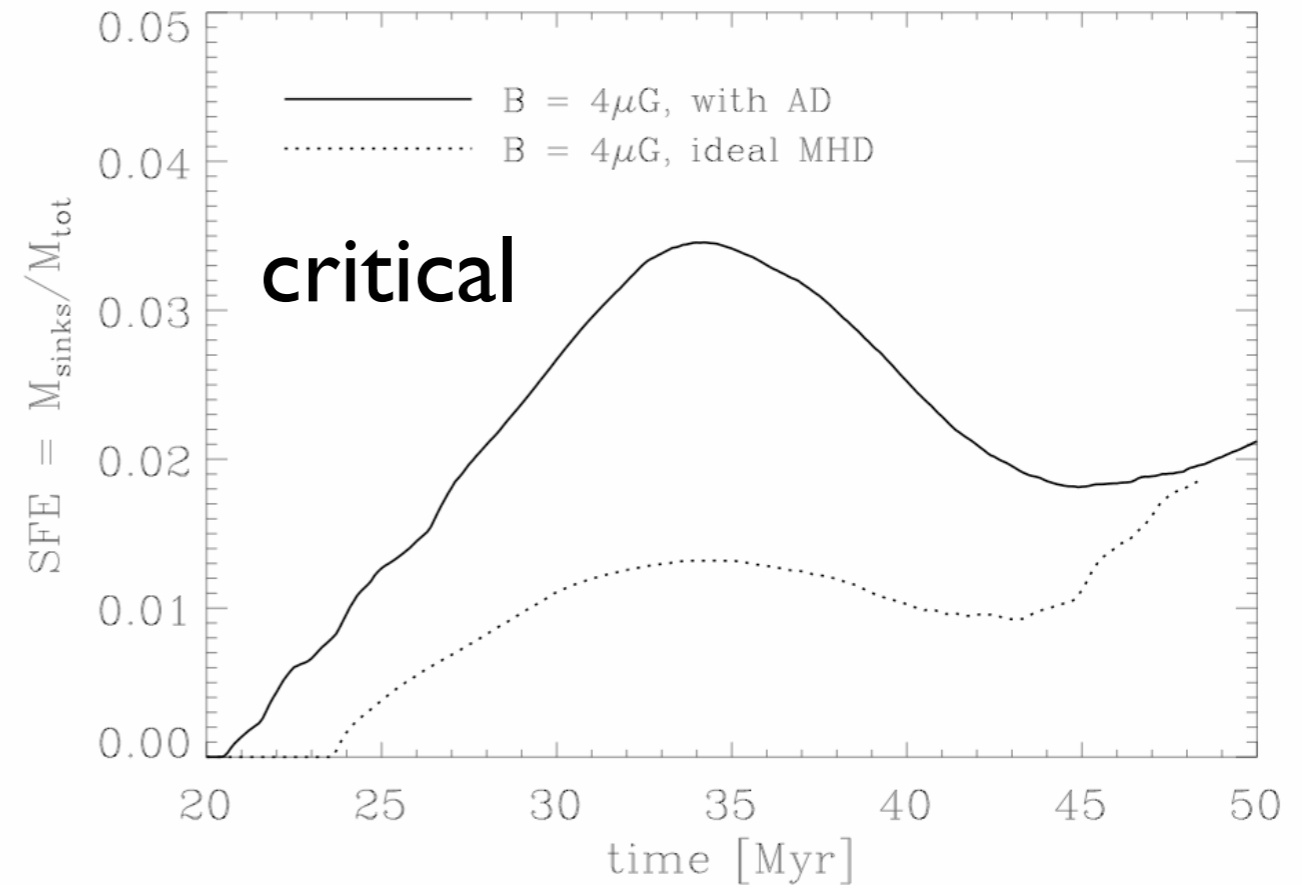
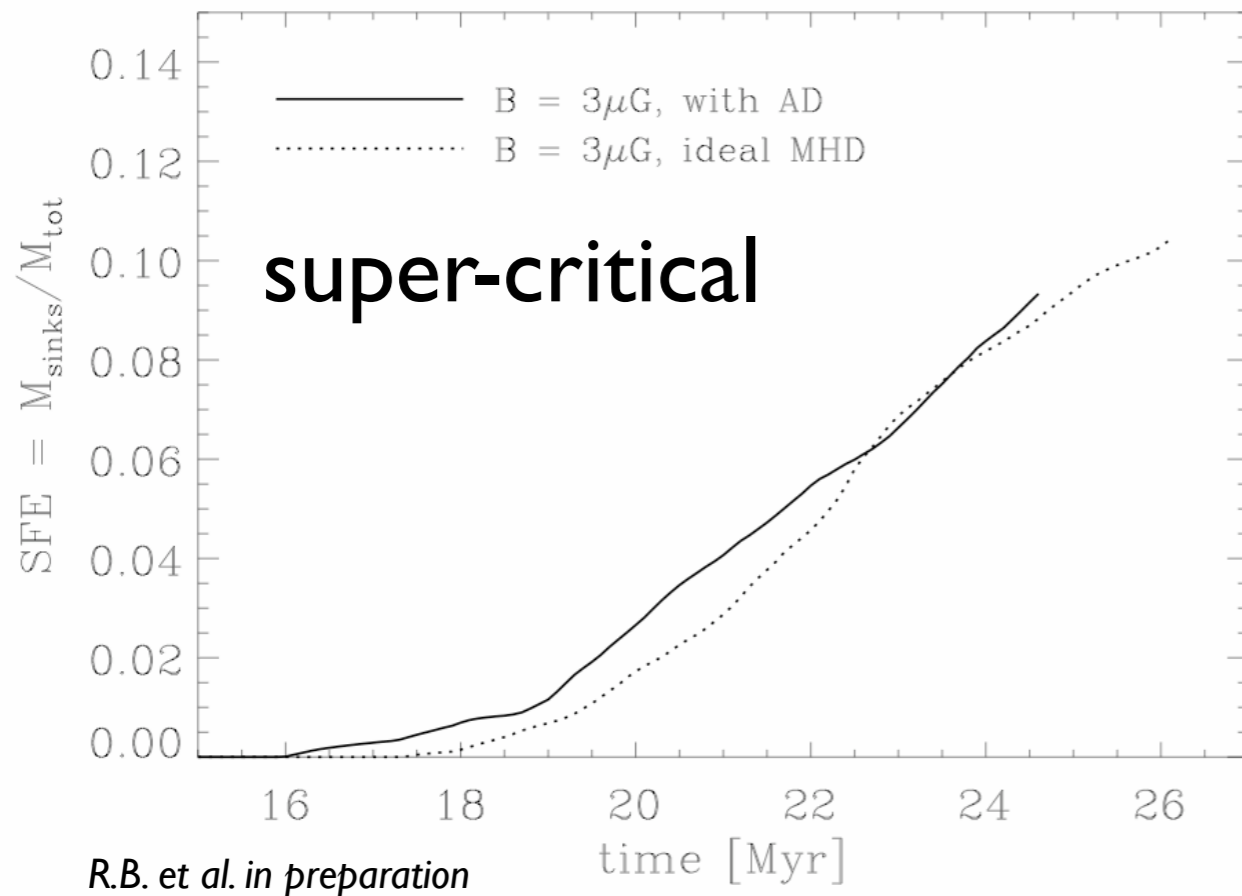


Boxsize 80.0 pc

with AD

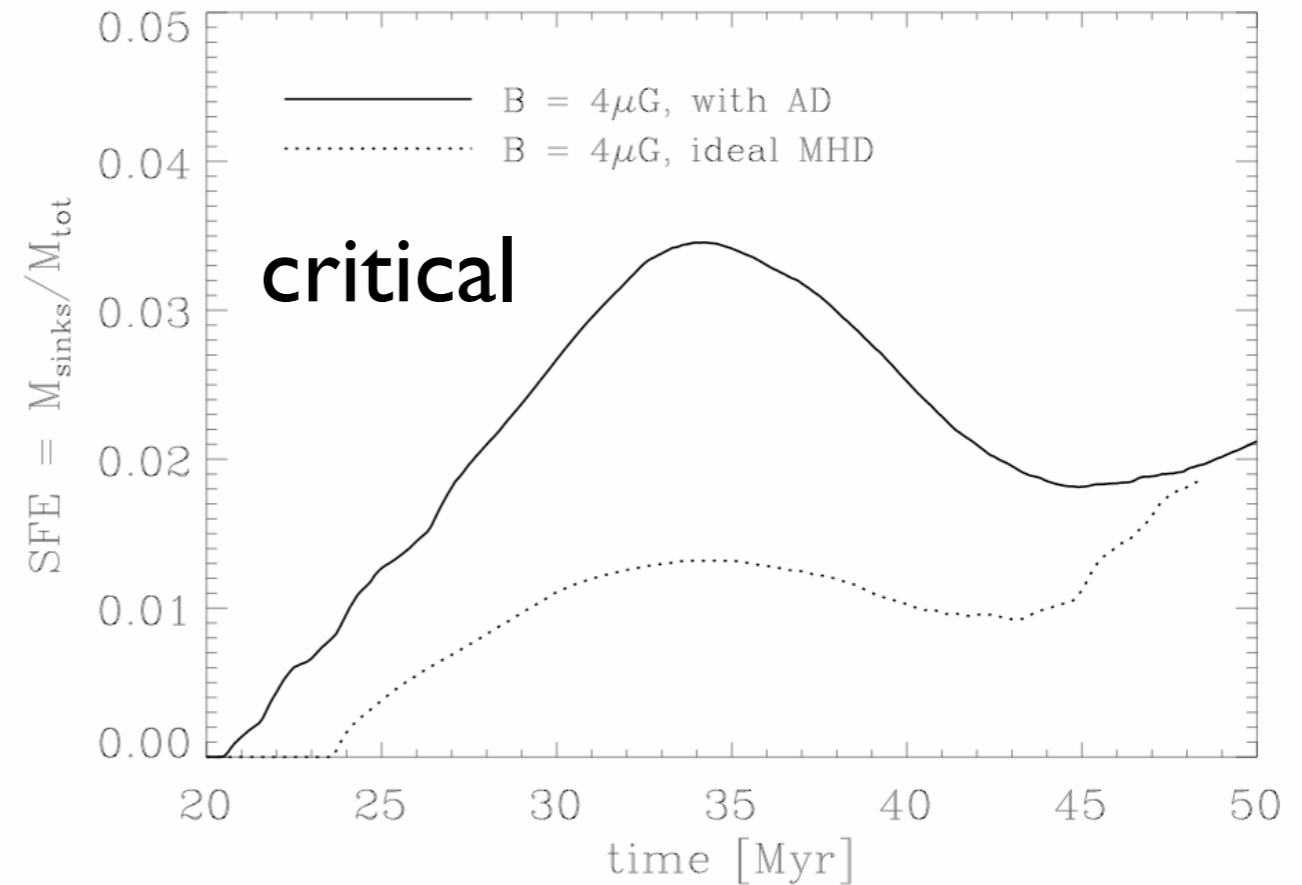
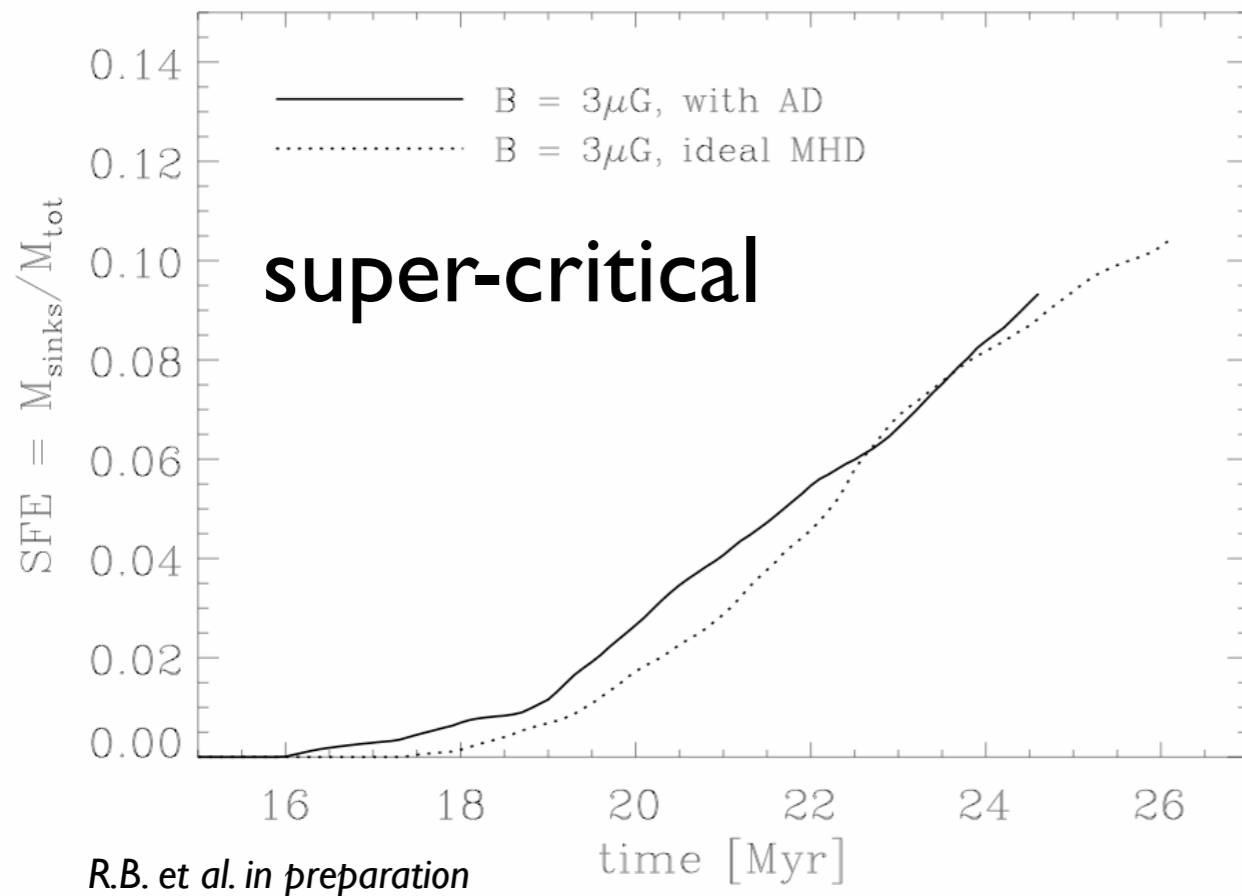
Formation of Molecular Clouds

Influence of Ambipolar Diffusion



Formation of Molecular Clouds

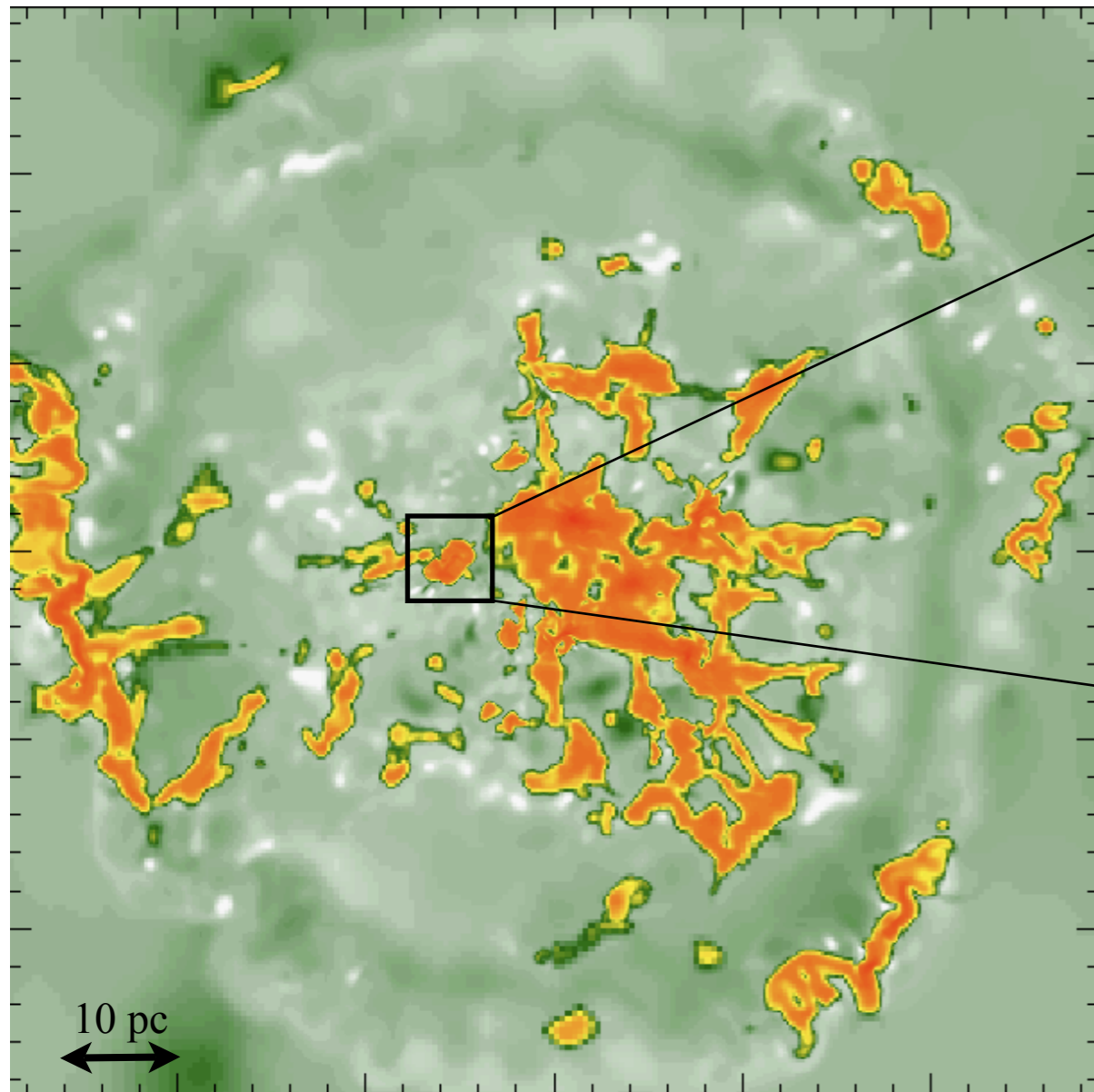
Influence of Ambipolar Diffusion



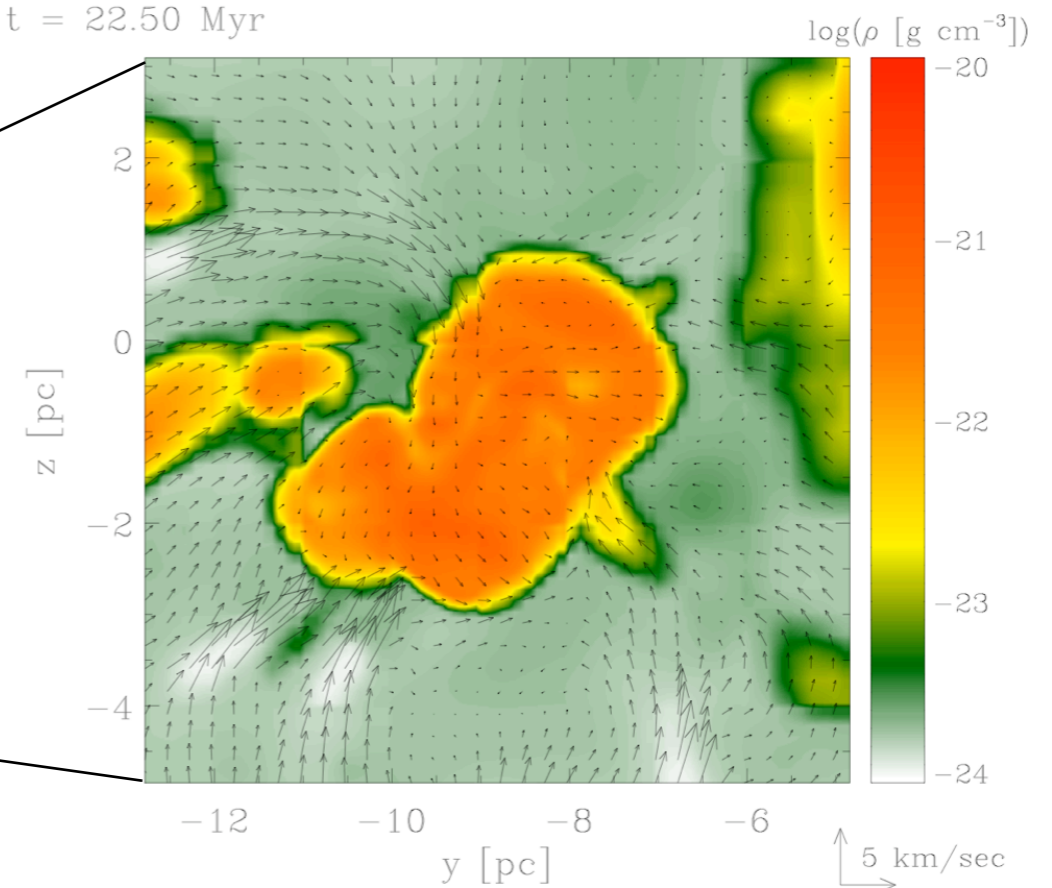
- Ambipolar diffusion is **not** a major player for star formation

Formation of Molecular Clouds

morphology and clump evolution



$t = 22.50$ Myr

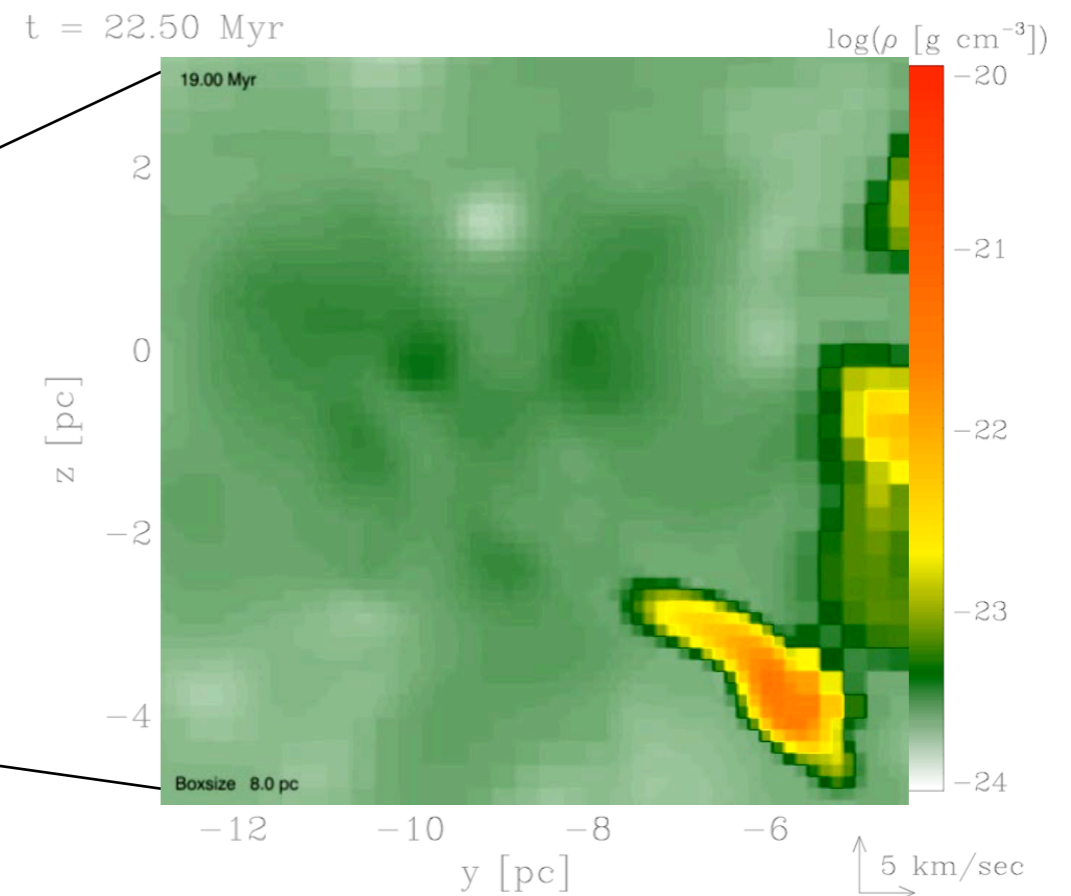
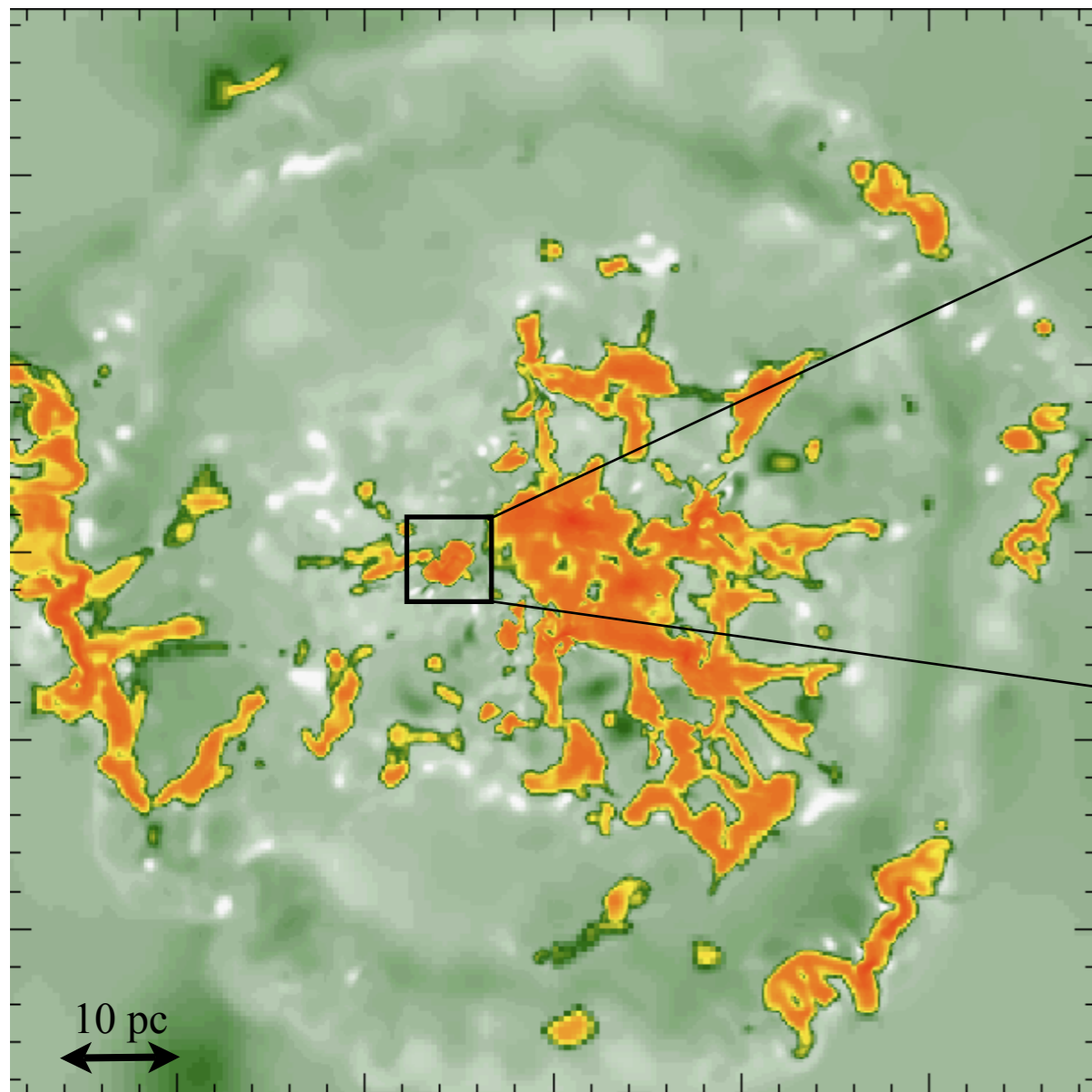


- MCs are inhomogeneous
- cold clumps embedded in warm atomic gas

- clumps growth by outward propagation of boundary layers and
- coalescence at later times

Formation of Molecular Clouds

morphology and clump evolution

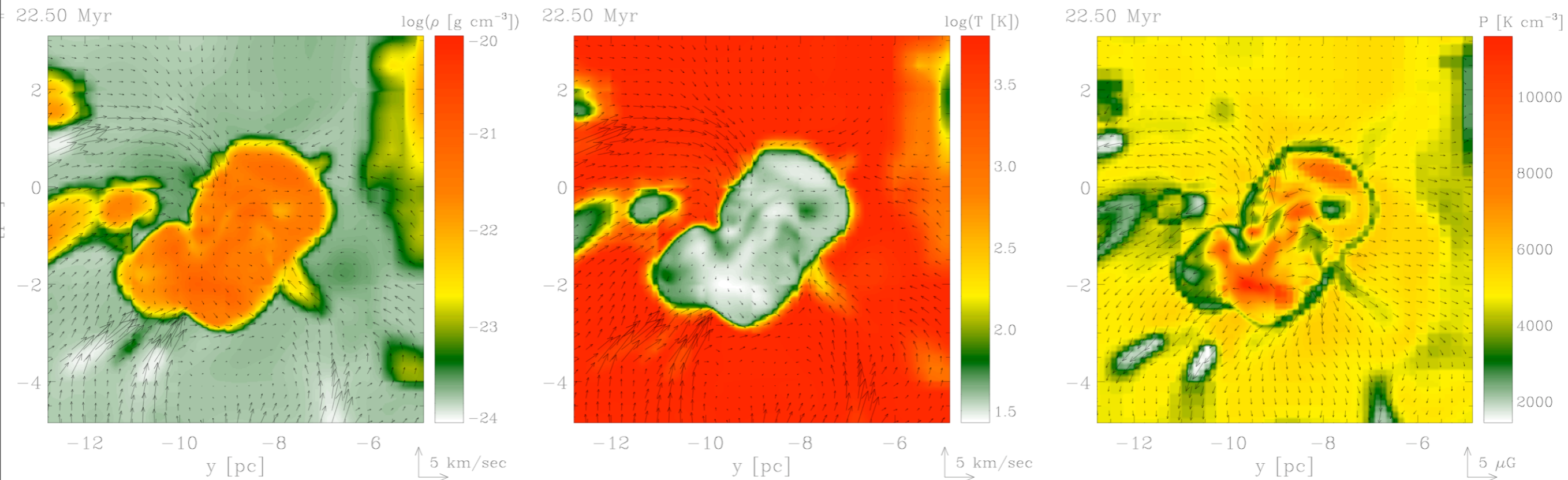


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Formation of Molecular Clouds

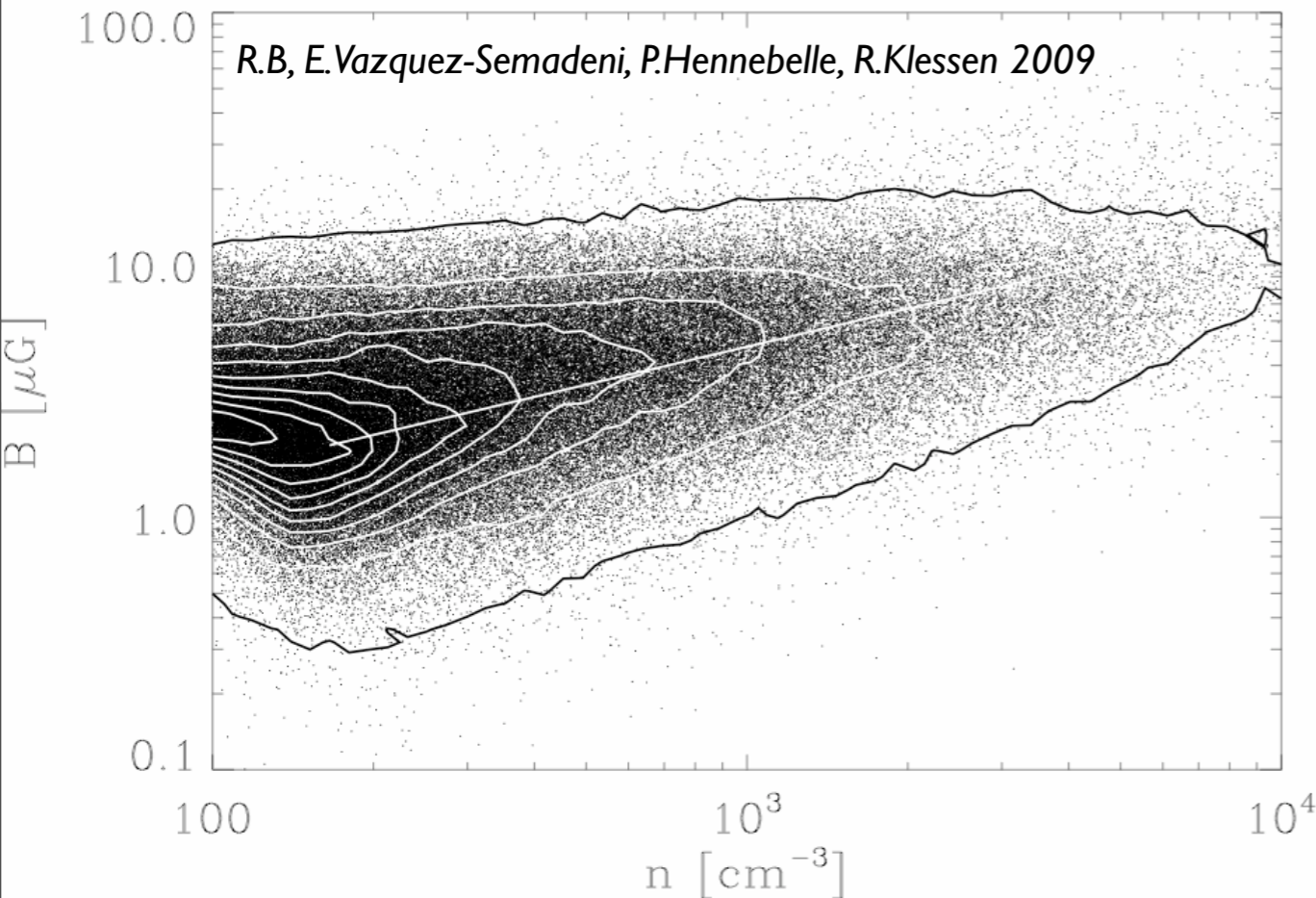
clump morphology



R.B, E.Vazquez-Semadeni, P.Hennebelle, R.Klessen 2009

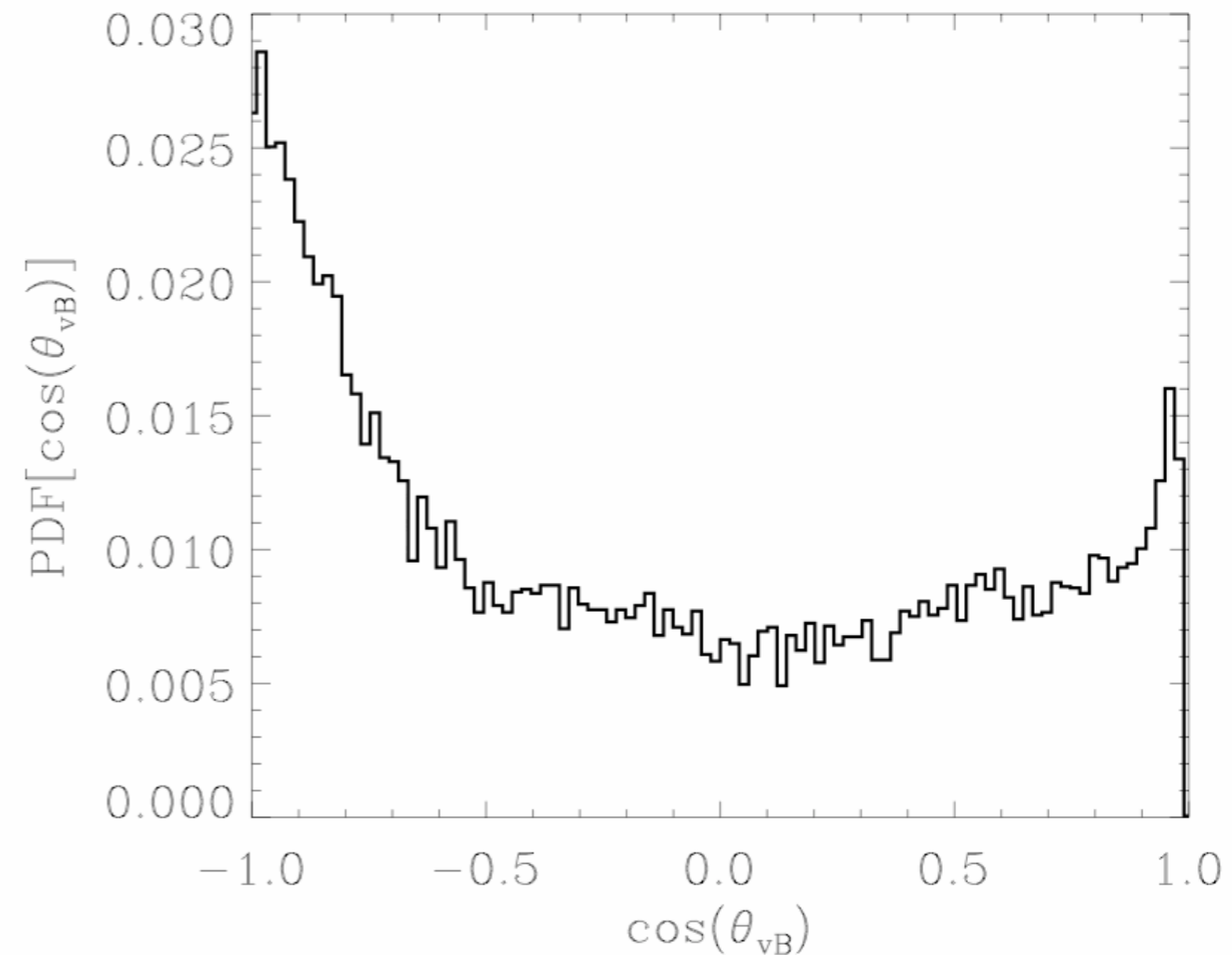
- cold clumps are in near pressure equilibrium (ram+thermal) with their warm surroundings
- in-falling gas streams along field lines

Formation of Molecular Clouds



- **large** scatter of magnetic field strengths: sub- and super-critical cores exist
- median slope: $B \propto n^{0.5}$ (e.g. *Crutcher 1999*)

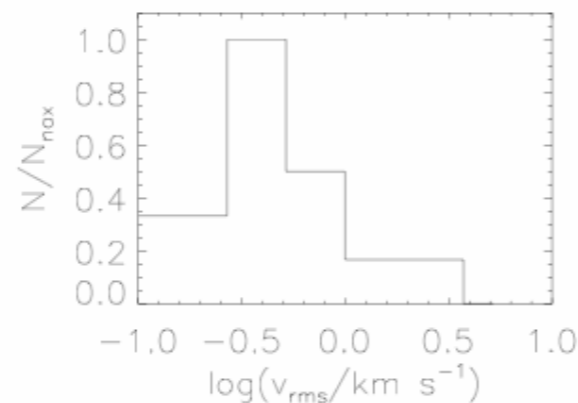
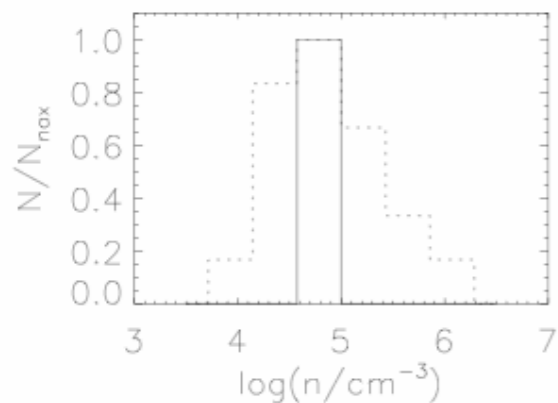
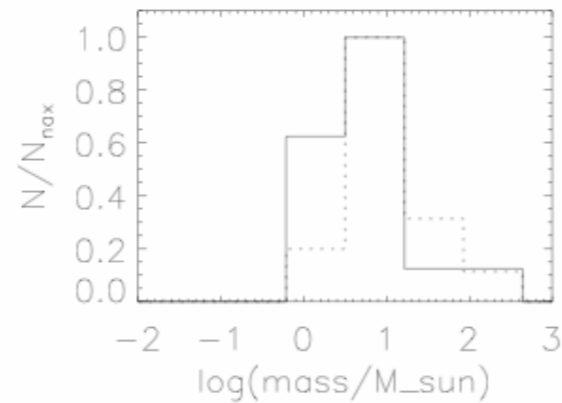
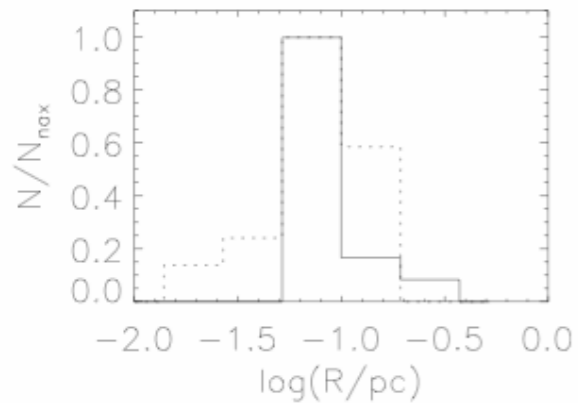
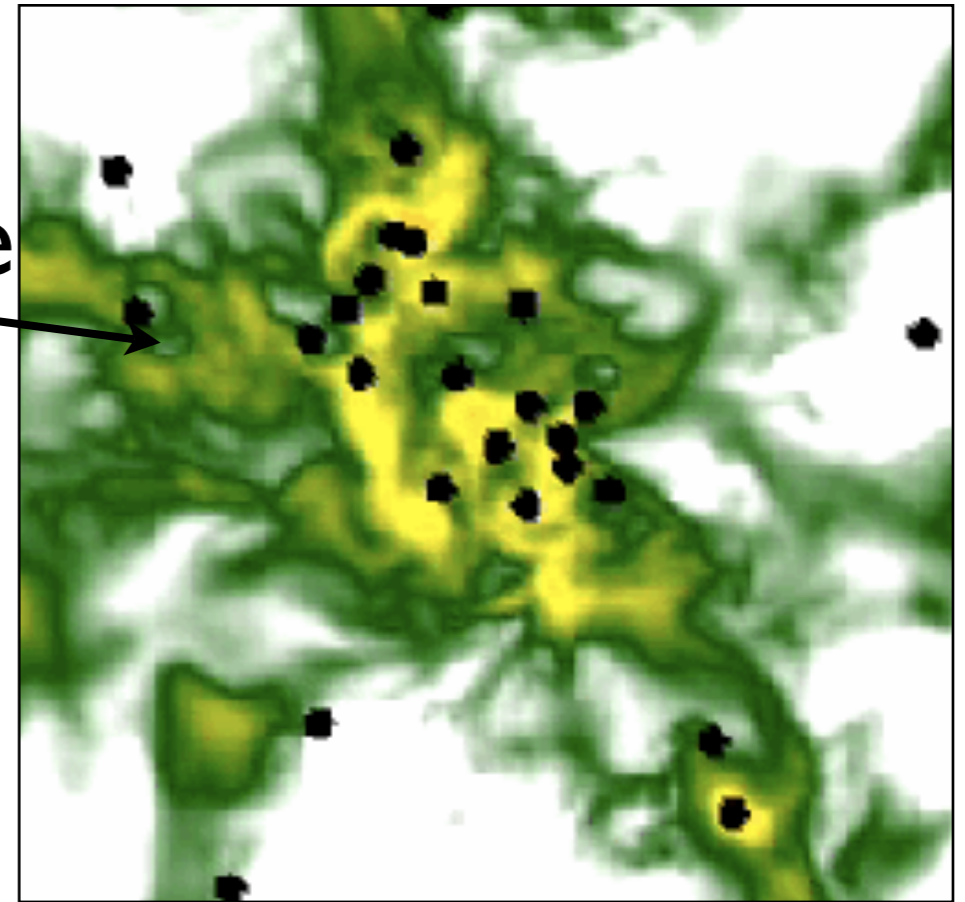
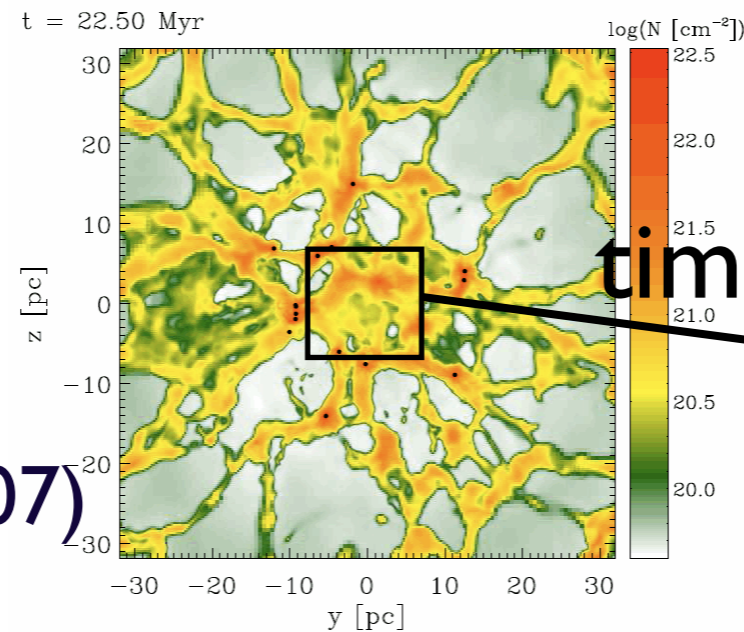
- **strong** correlation of gas streams and magnetic field lines



Formation of Molecular Clouds

global contraction phase

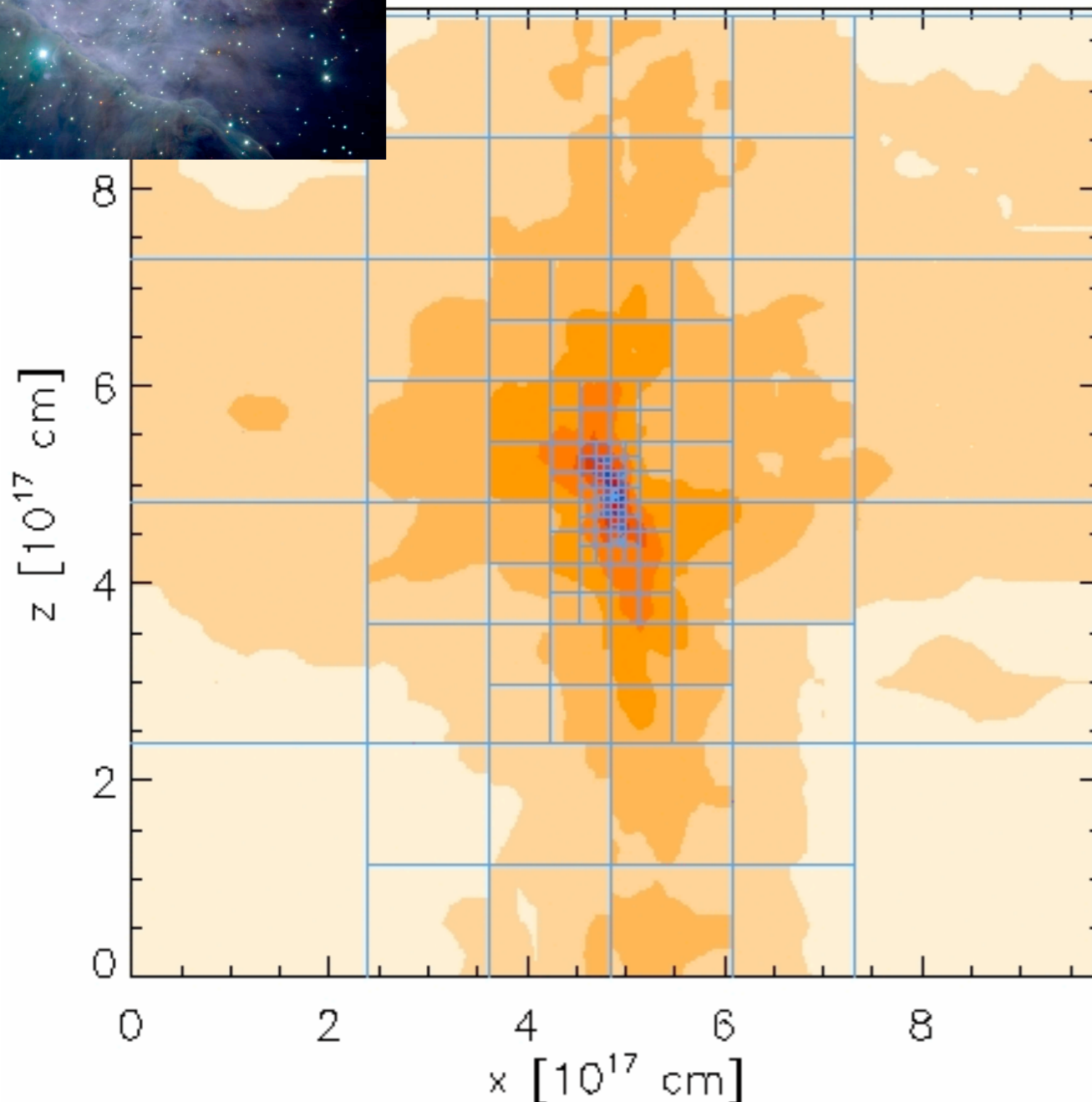
centre of the cloud
→ birthplace for
massive stars?
(eg. Zinnecker & Yorke 2007)



comparison of core properties
with observation of Cygnus X
by *Motte et al 2007*

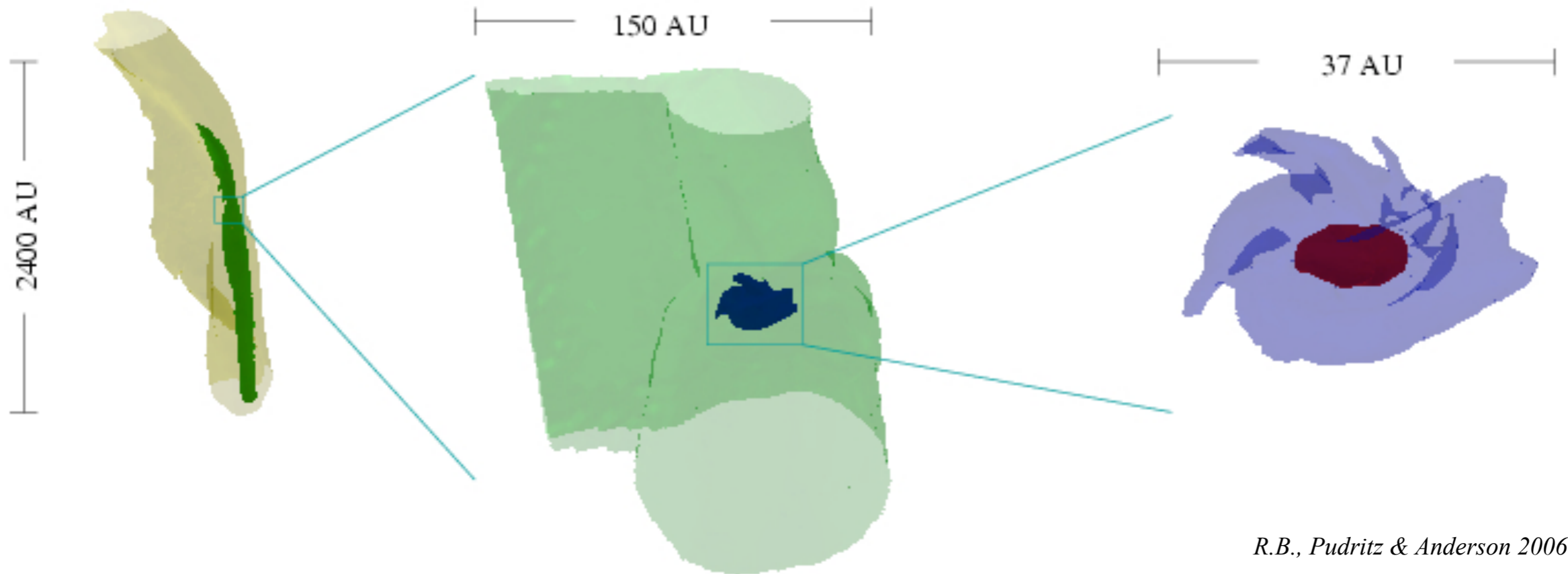
Collapse of turbulent cores

SF out of **turbulent**
molecular clouds
(e.g. *Mac Low & Klessen 2004*)



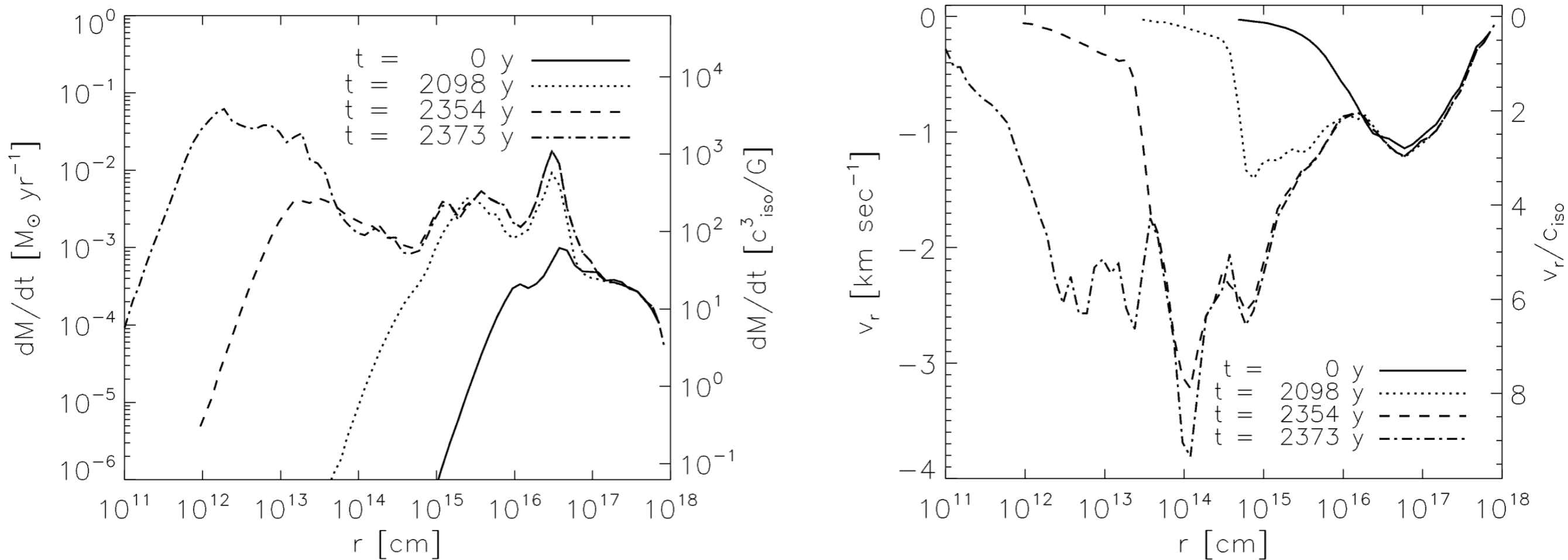
- Initial data from *Tilley & Pudritz 2004*: **ZEUS** simulations of core formation within a supersonic **turbulent** environment
- $L = 0.32 \text{ pc}$, $M_{\text{tot}} = 105 M_{\text{sol}}$
- Follow the collapse of the densest most **massive** region: $\sim 23 M_{\text{sol}}$
- Final resolution: $\sim R_{\text{sol}}$ (27 refinement levels)

Collapse of turbulent cores



- **Filament** with an attached sheet
- small **disk** within the filament (perpendicular)
- adiabatic (optically thick) core
- very efficient gas **accretion** through the filament

Formation of massive stars

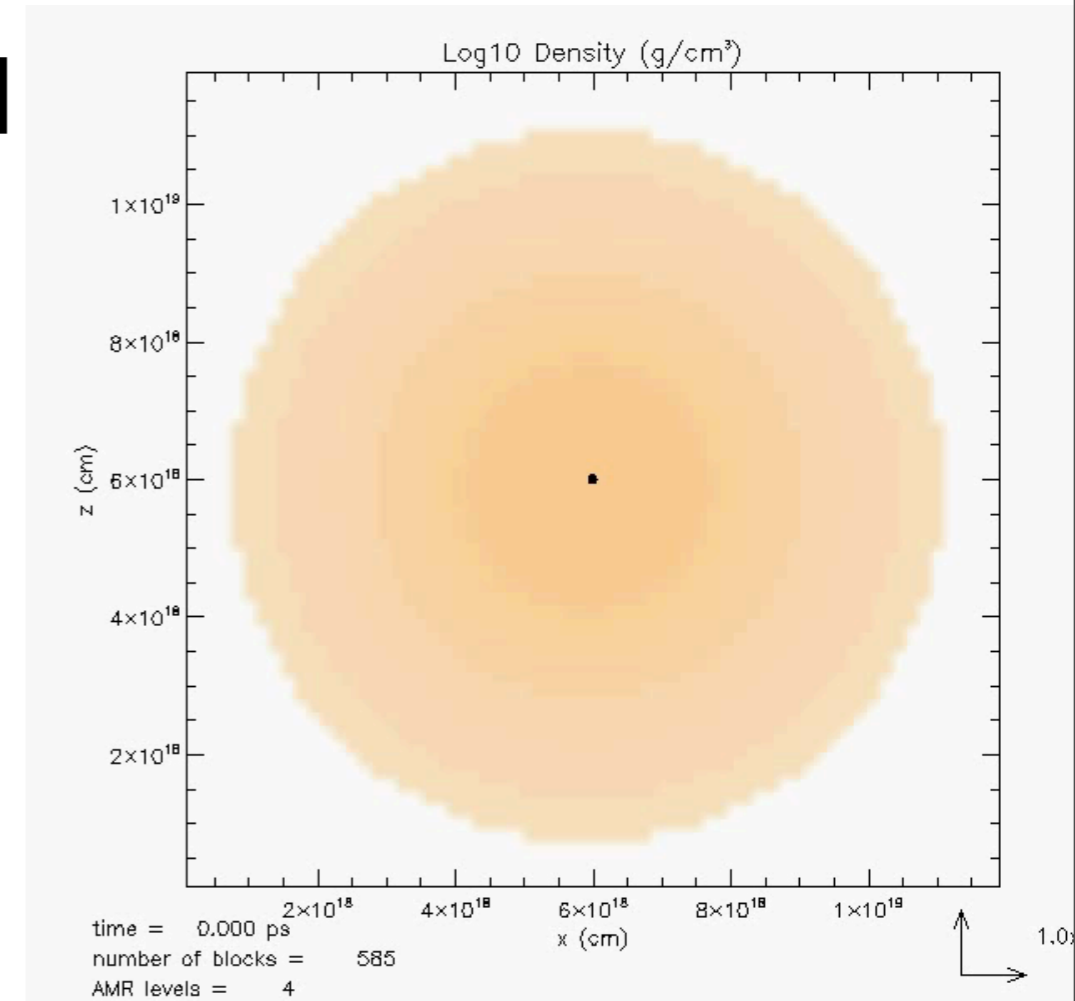


- Very **high** accretion rates through dense filaments:
up to $10^{-3} - 10^{-2} M_{\text{sol}}/\text{year}$
- Mass accretion rates are higher than limits from radiation pressure of **massive** stars (e.g. *Wolfire & Cassinelli 1987*: $10^{-3} M_{\text{sol}}/\text{year}$)

Massive Star Formation: Dynamics of HII Regions

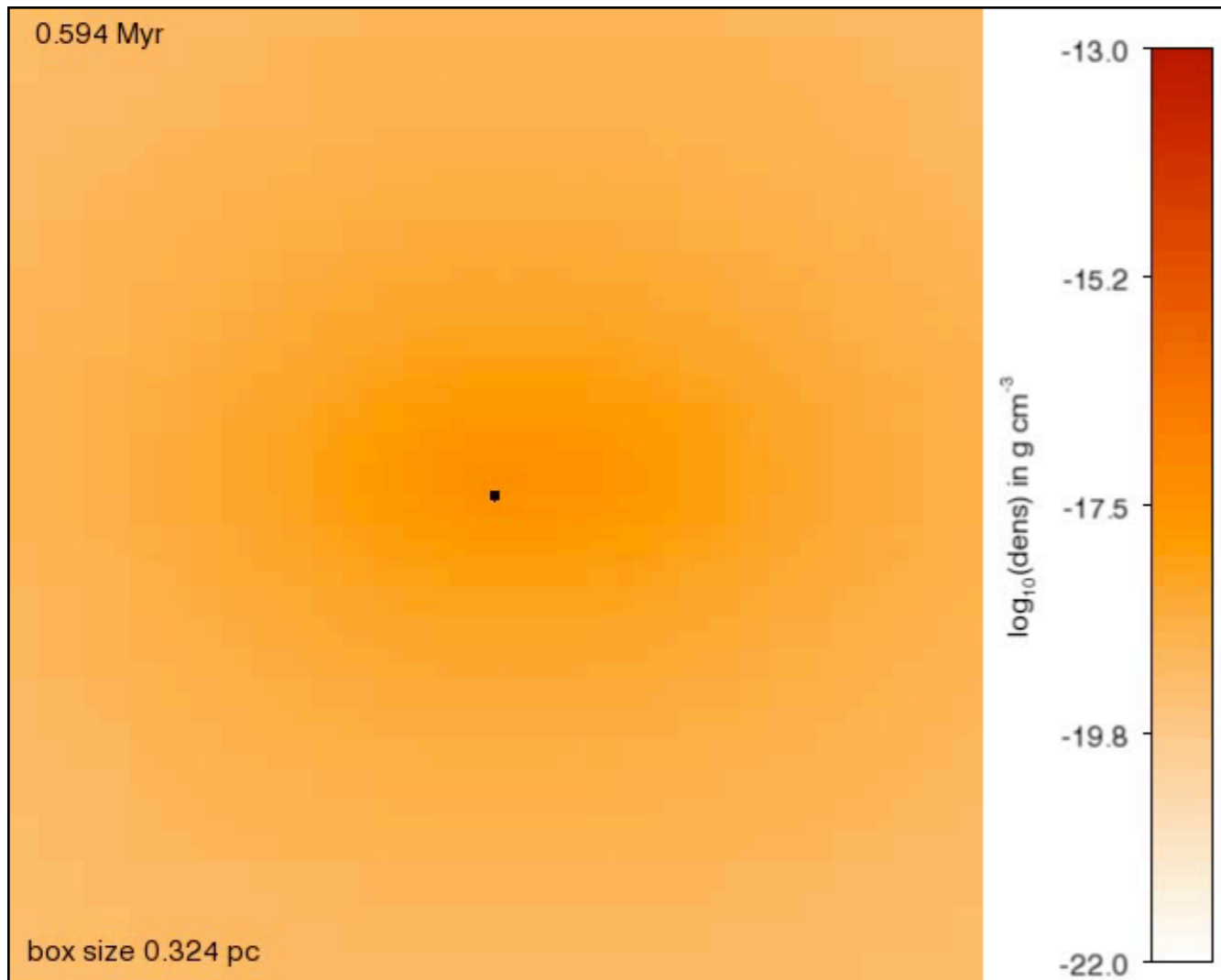
3D Simulations of collapsing cloud cores with **ionization feedback** from young massive stars

- massive core with $M_{\text{core}} = 1000 M_{\text{sol}}$
- flat core with $r = 0.5 \text{ pc}$ and $\rho \sim r^{-1.5}$
- initial core rotation with $\beta = 0.05$
- accreting sink particles \Rightarrow luminosity and temperature using ZAMS (*Paxton 2004*)
- highest grid resolution $\sim 100 \text{ AU}$

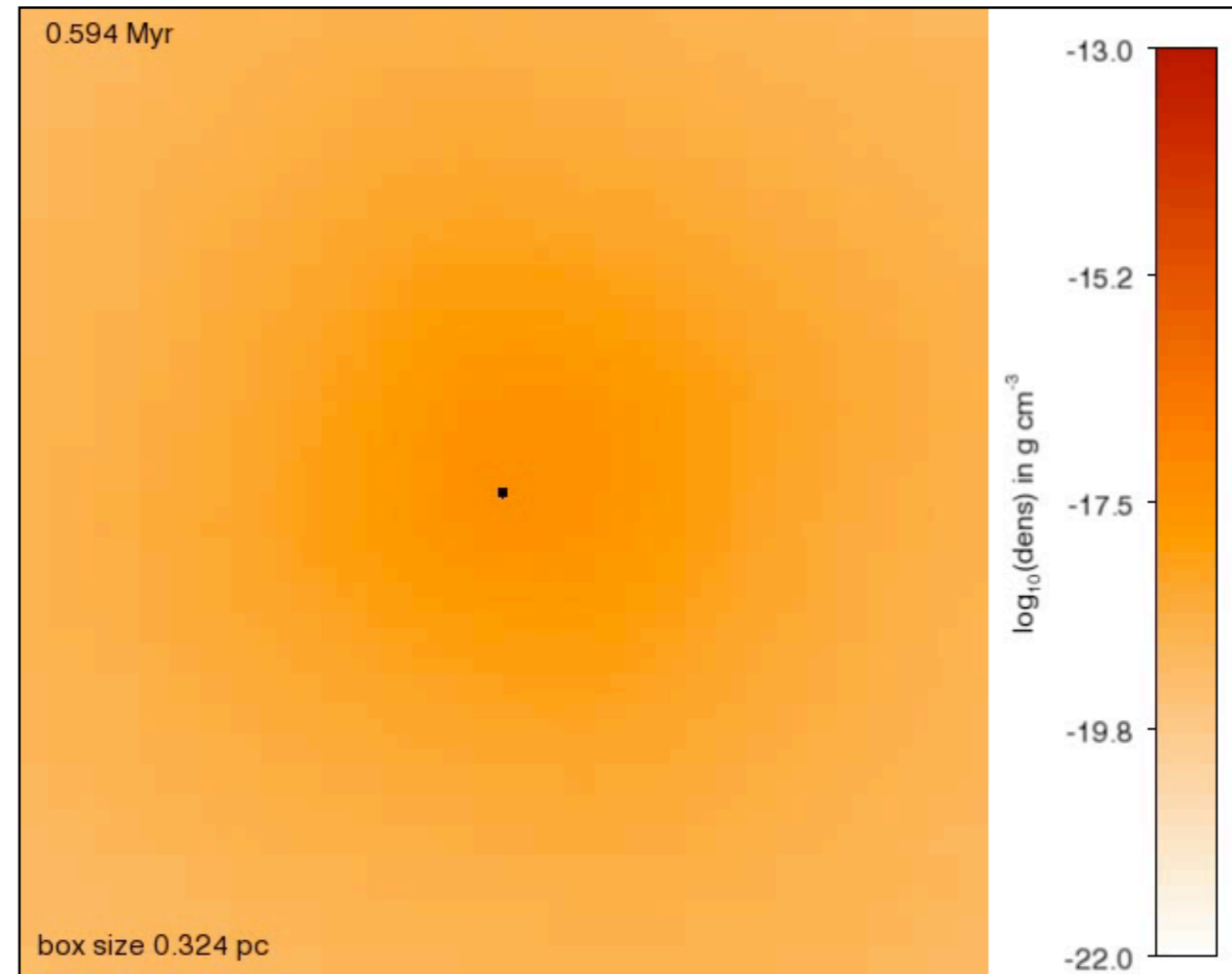


Massive Star Formation: Dynamics of HII Regions

Simulations by Thomas Peters (ITA)



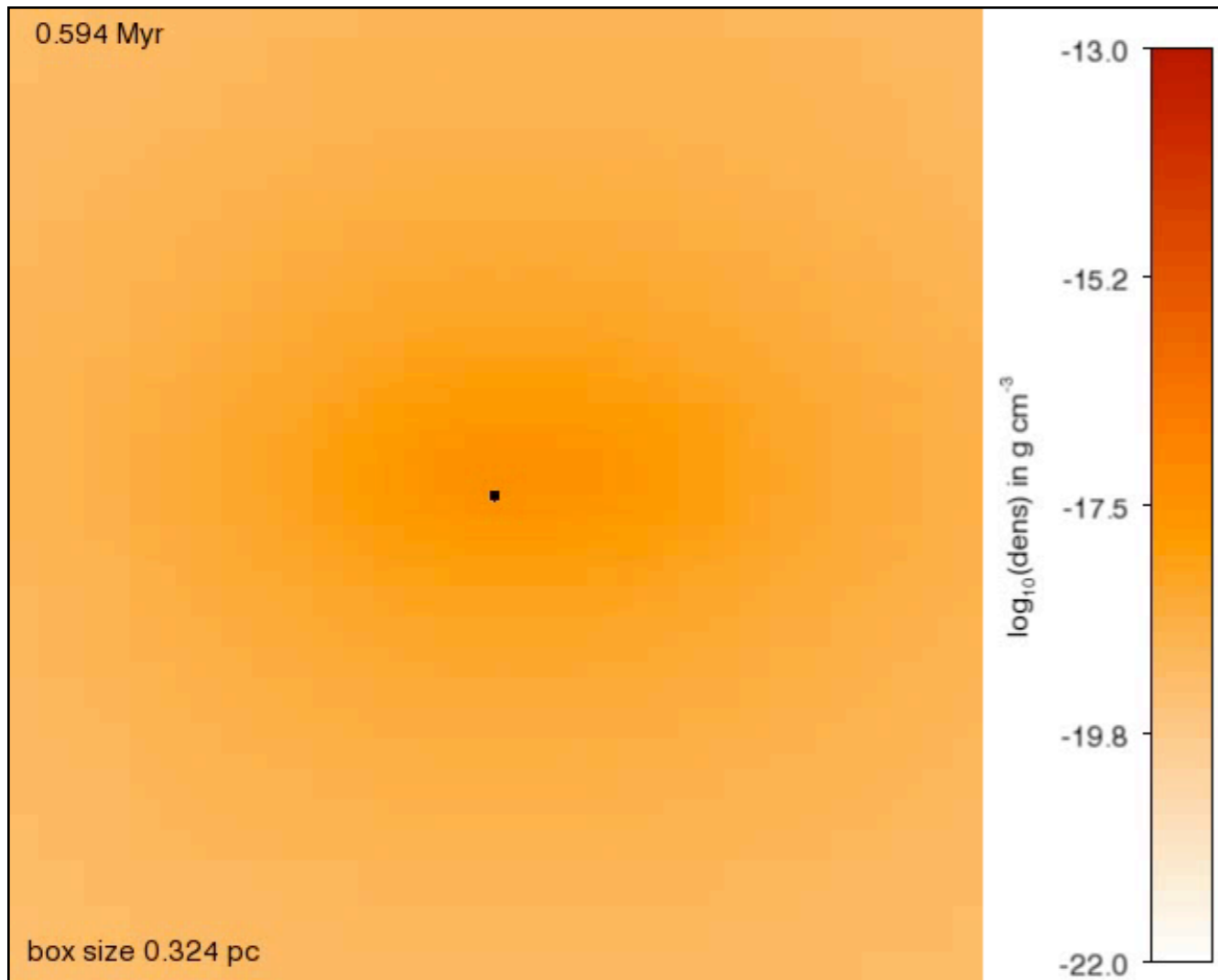
Disk edge on



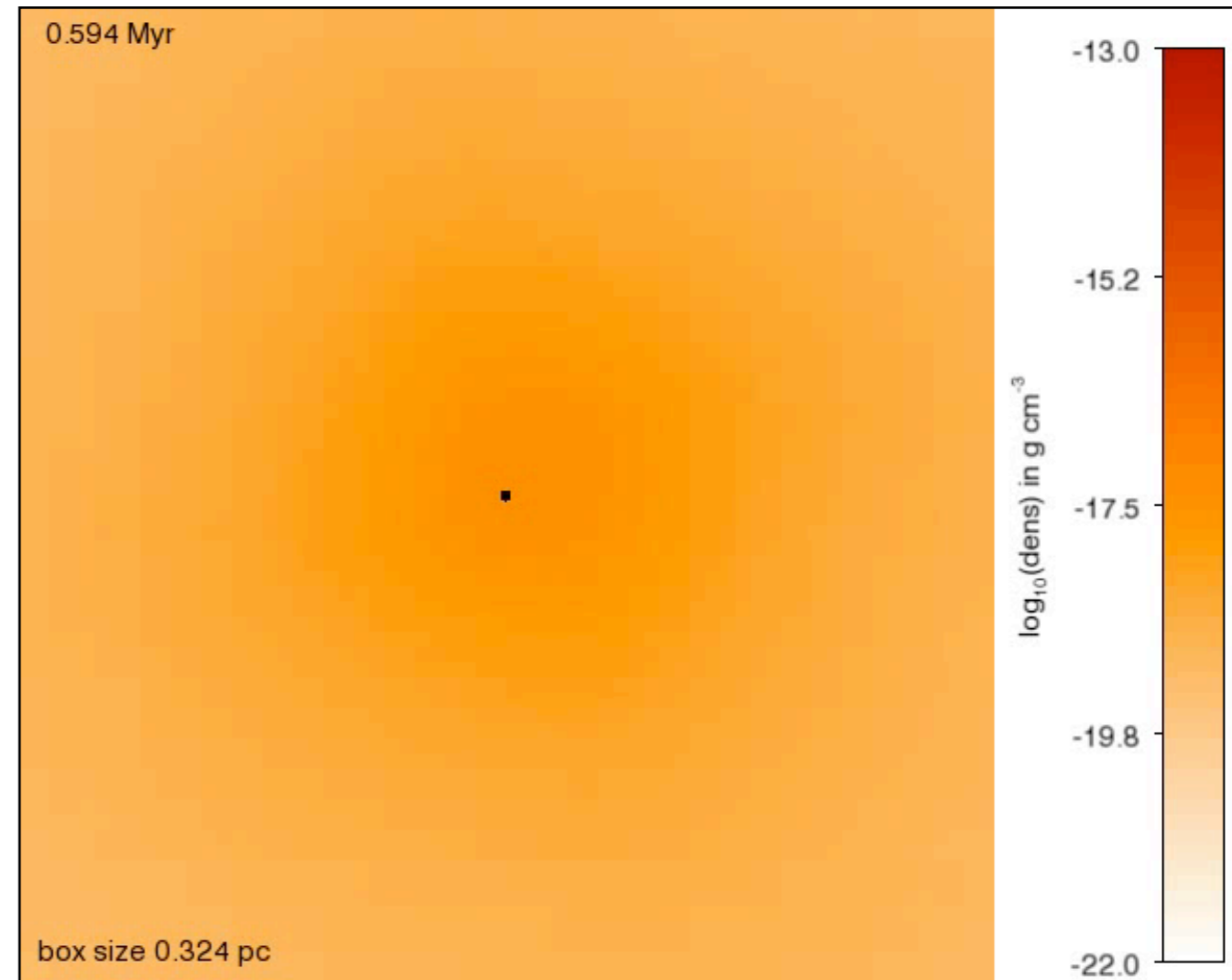
Disk plane

Massive Star Formation: Dynamics of HII Regions

Simulations by Thomas Peters (ITA)

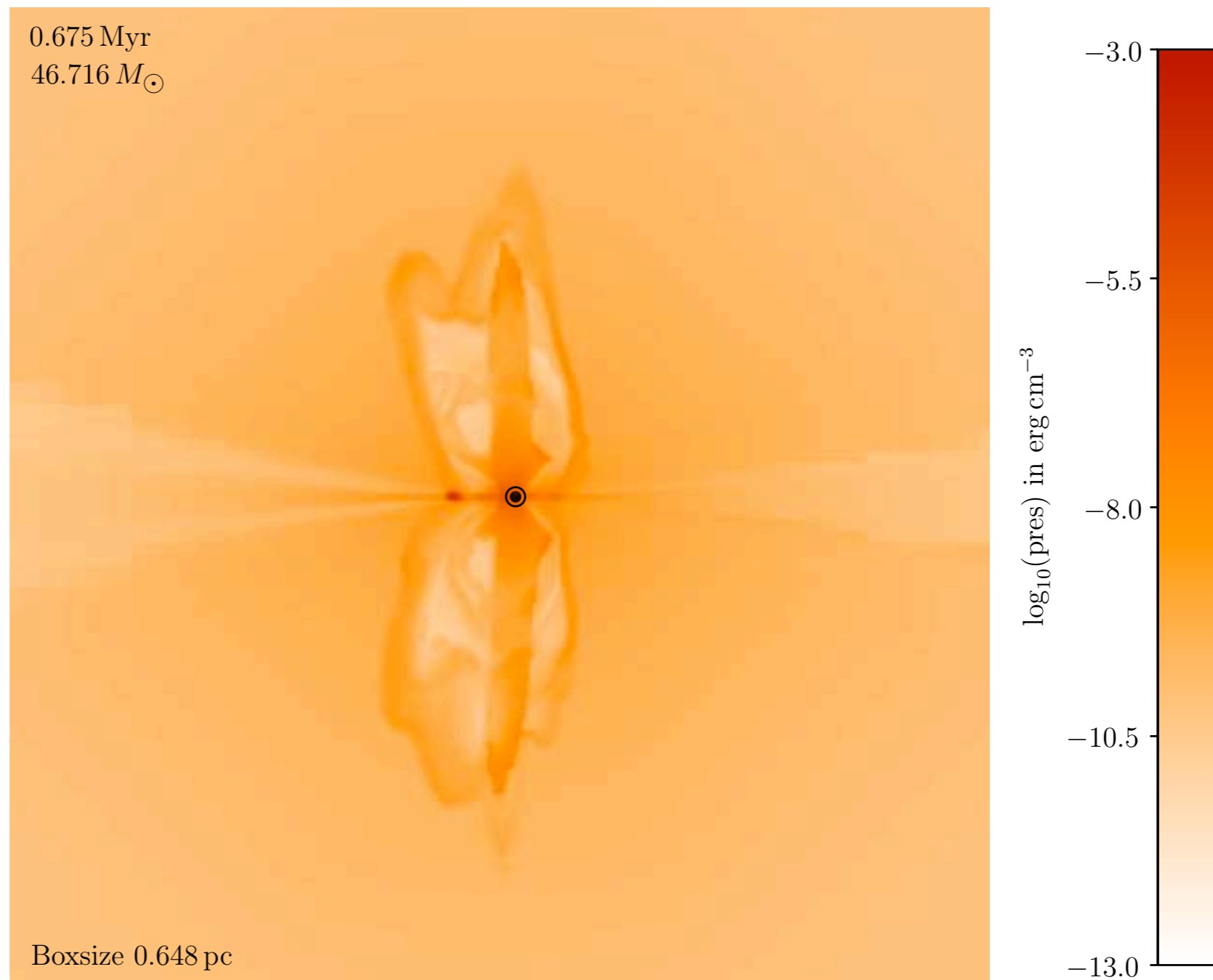


Disk edge on



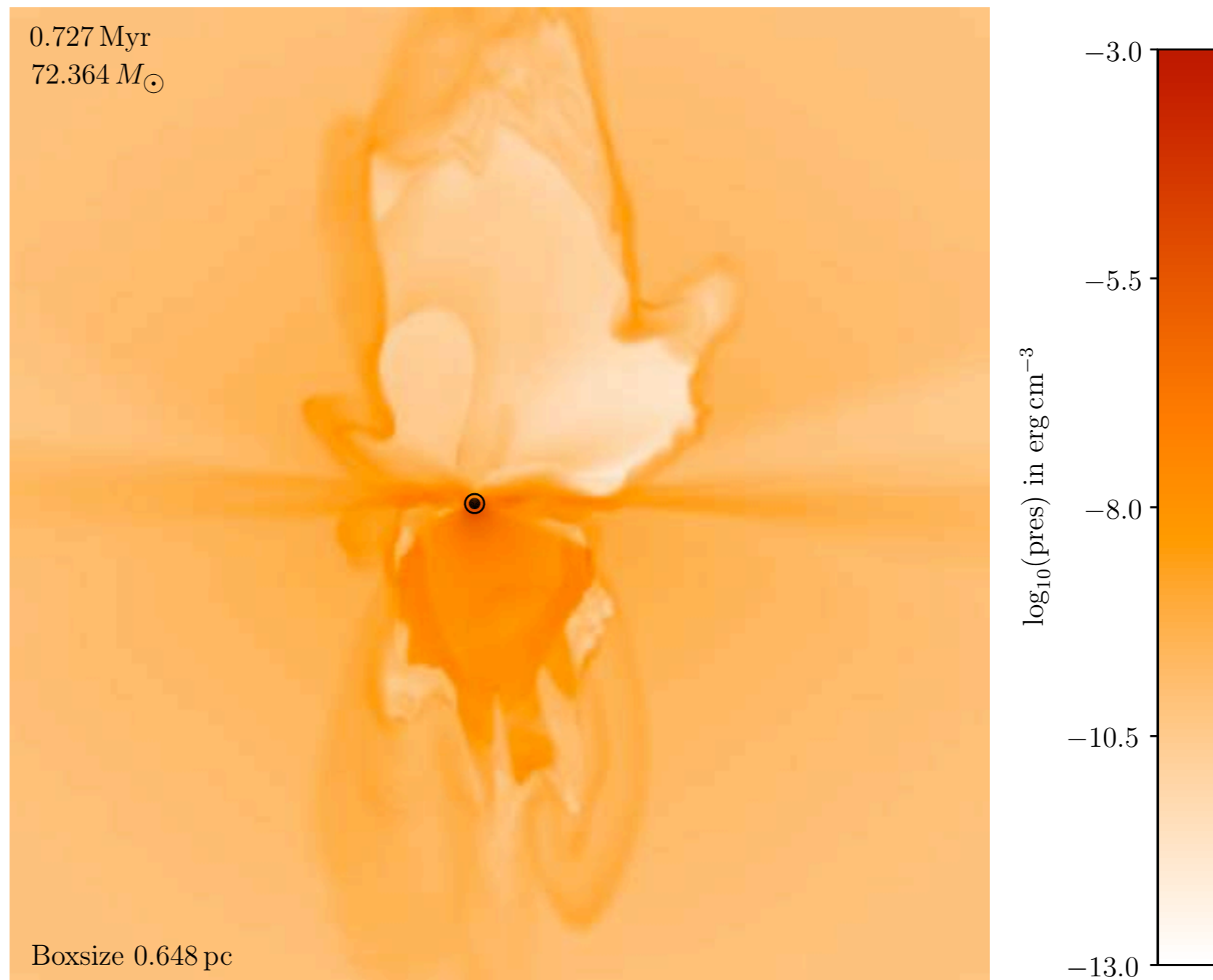
Disk plane

Dynamics of the H II Region and Outflow



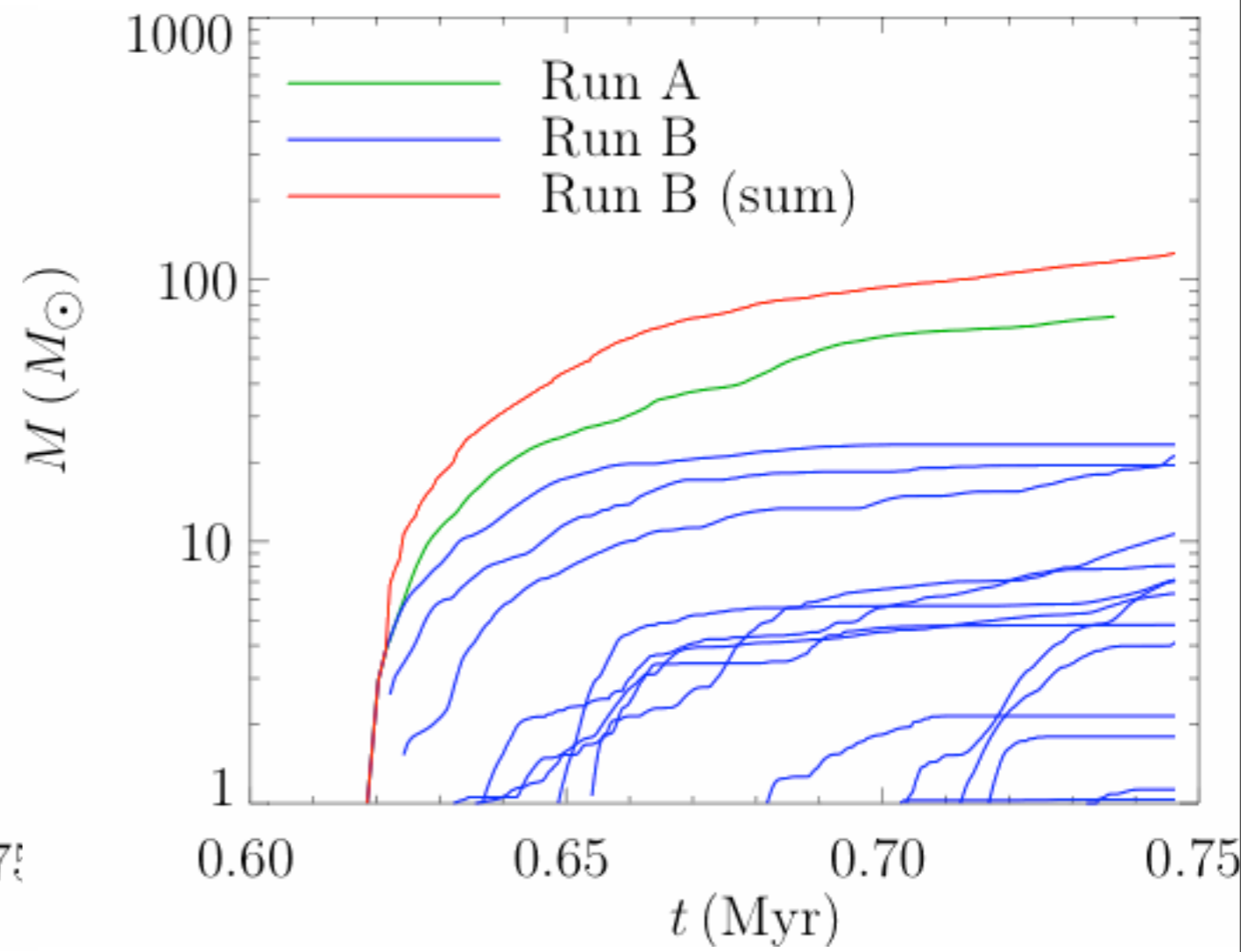
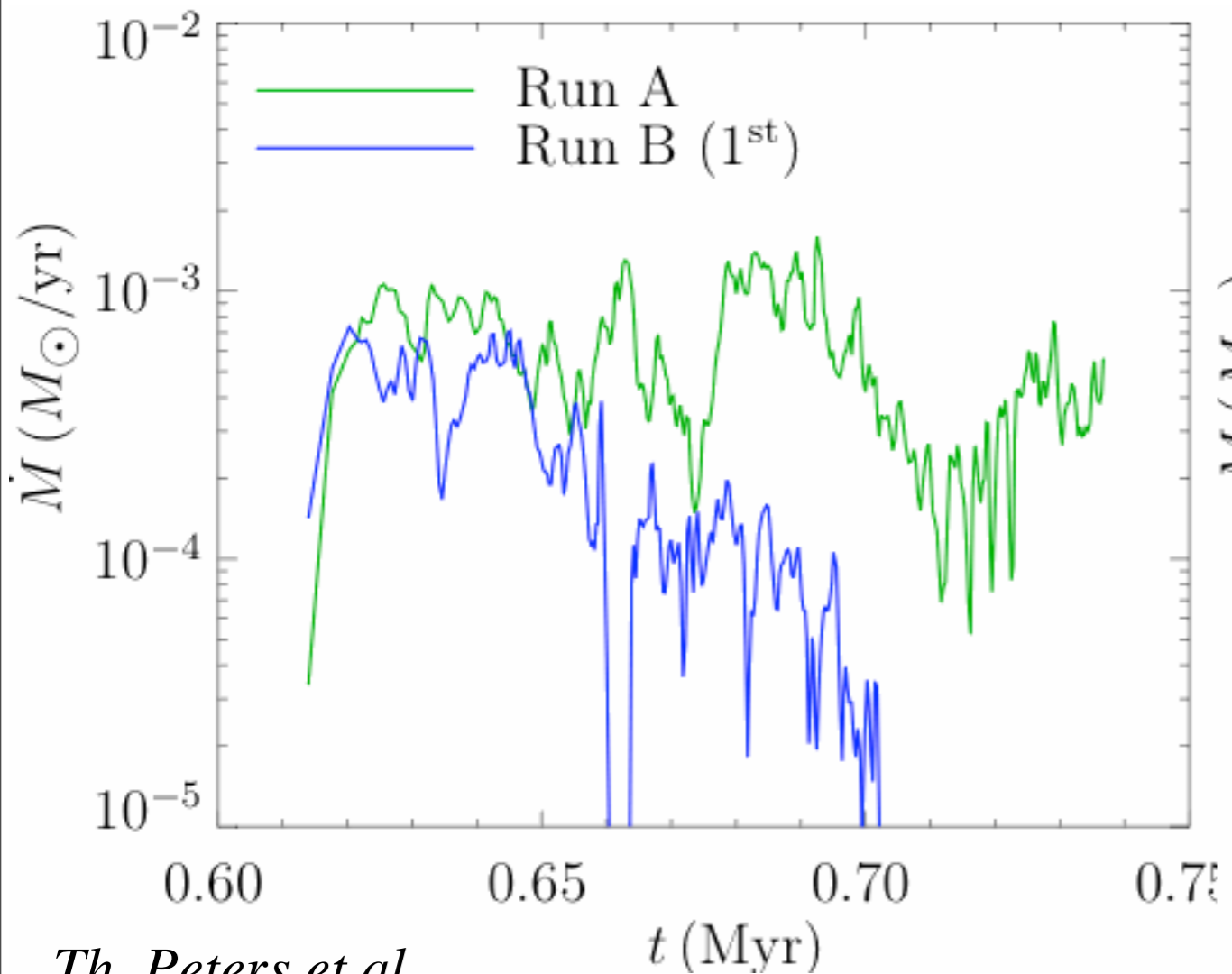
- ionization drives bipolar outflow
- pressure-driven expansion of shell
- thin-shell instability leads to fingers

Dynamics of the H II Region and Outflow



- size and morphology of H II region is highly variable
- cometary H II region totally reverses within less than 10 kyr
- changes like this have been observed!

Multiple protostars: Dynamics of the H II Region

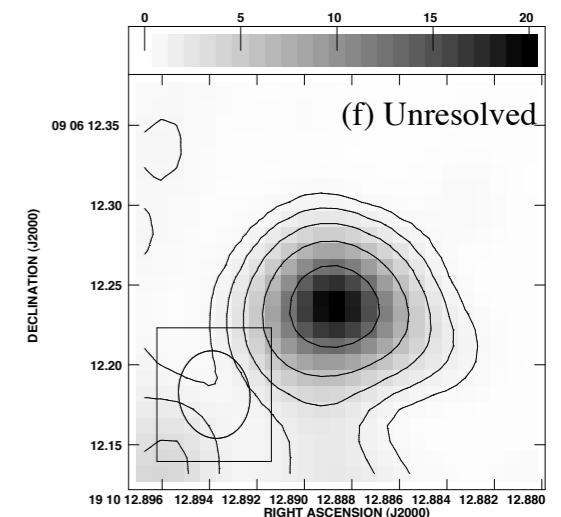
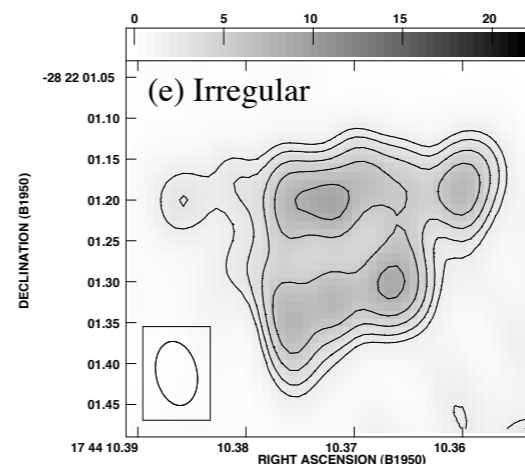
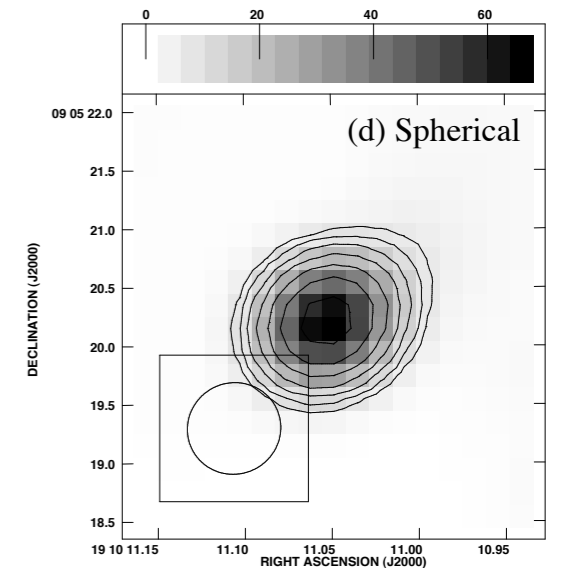
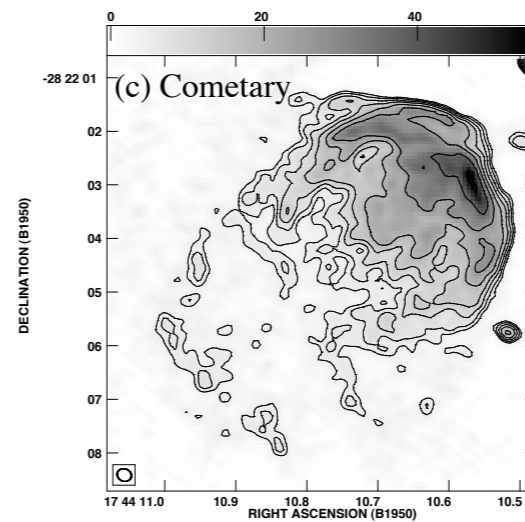
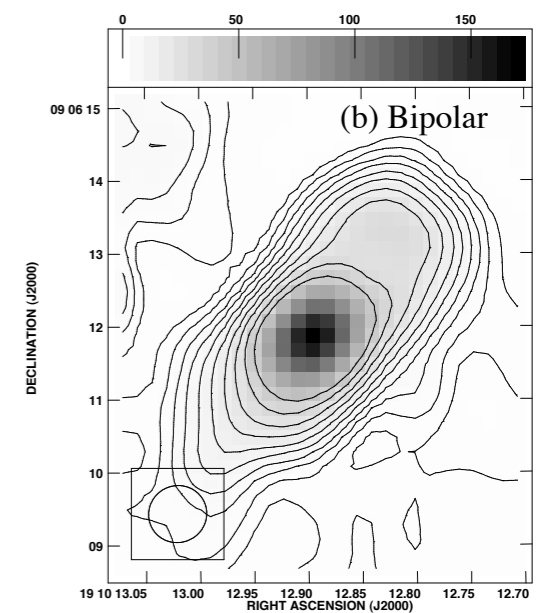
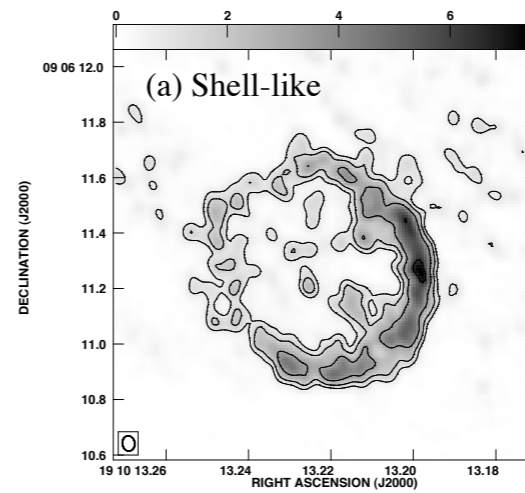


Th. Peters et al.

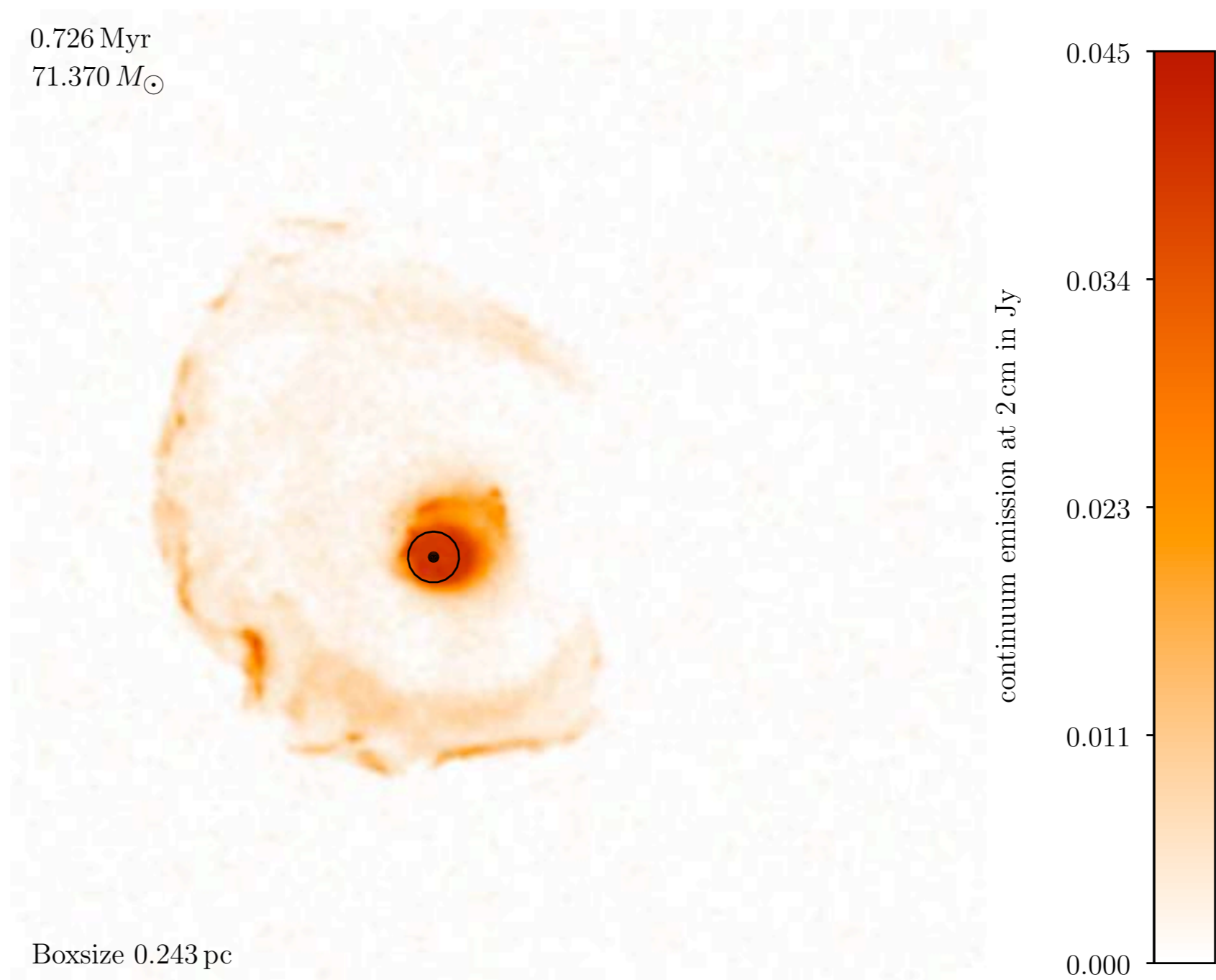
- ionization feedback does **not** shut off accretion
- **fragmentation**-induced starvation
- massive stars form in cluster

Classification of UC H II Regions

- comparison with De Pree et al. 2005 classification of UC H II regions in W49A and Sagittarius B2
- “irregular” is any resolved region which does fall into one of the other categories

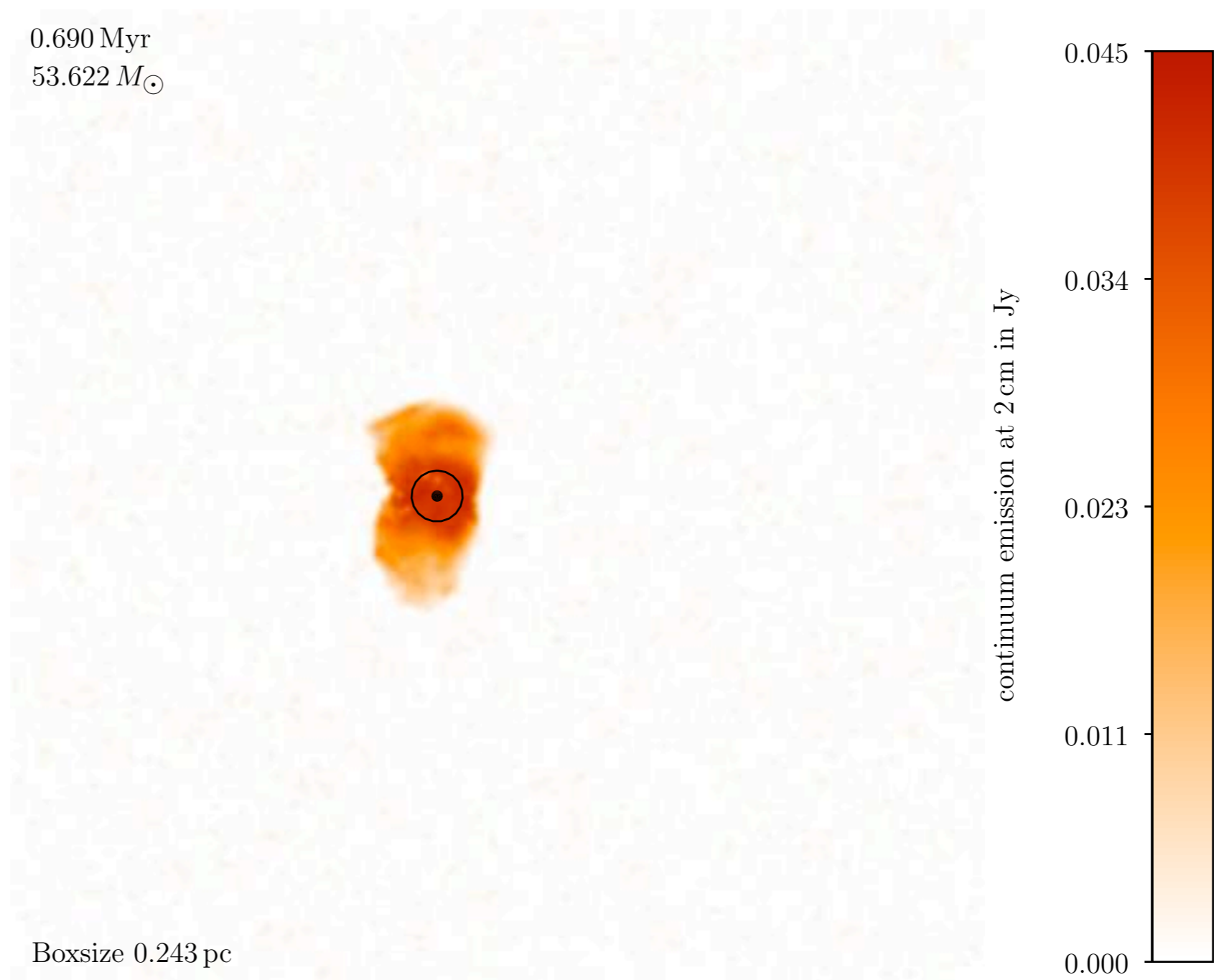


H II Region Morphologies



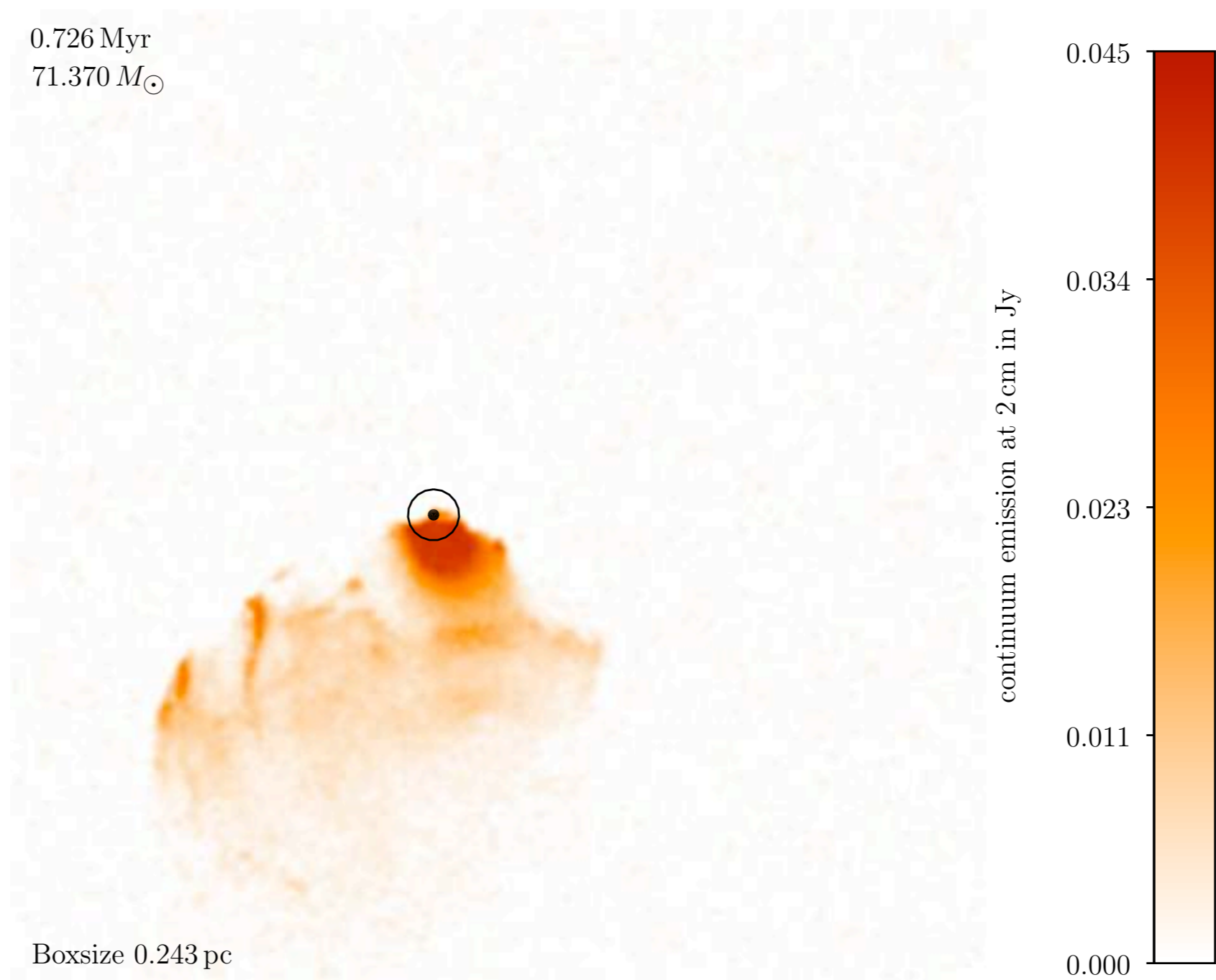
shell-like morphology

H II Region Morphologies



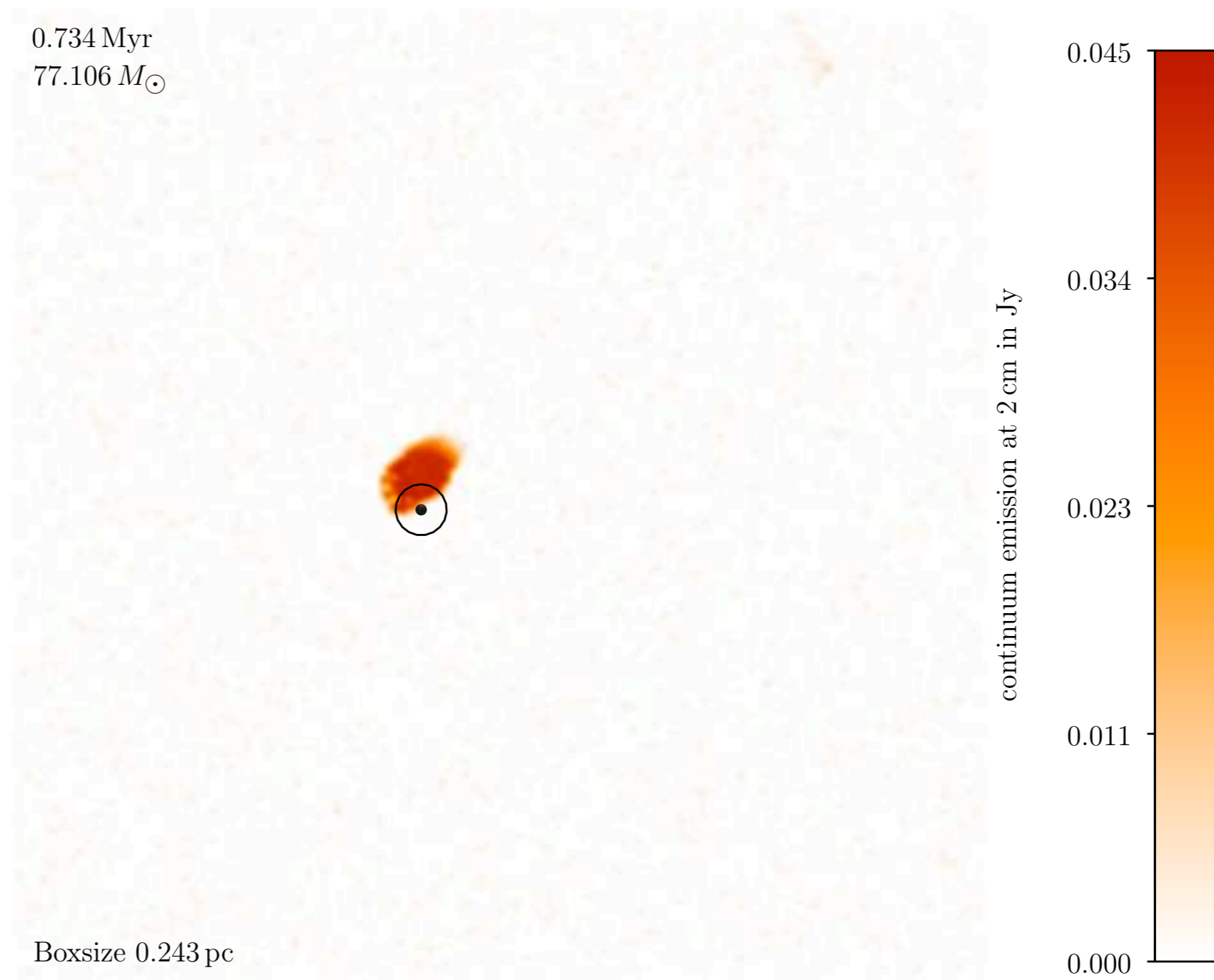
bipolar morphology

H II Region Morphologies

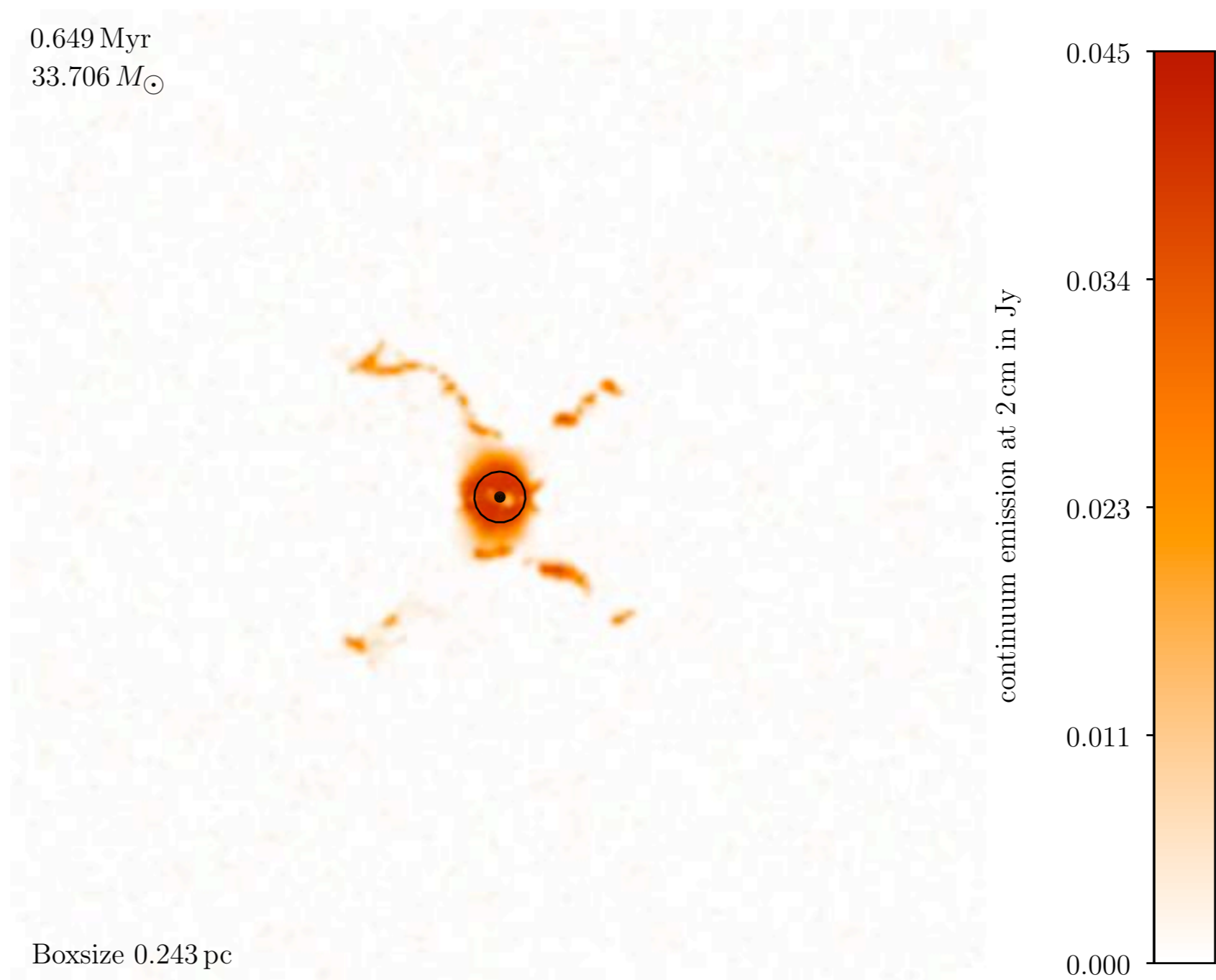


cometary morphology

H II Region Morphologies

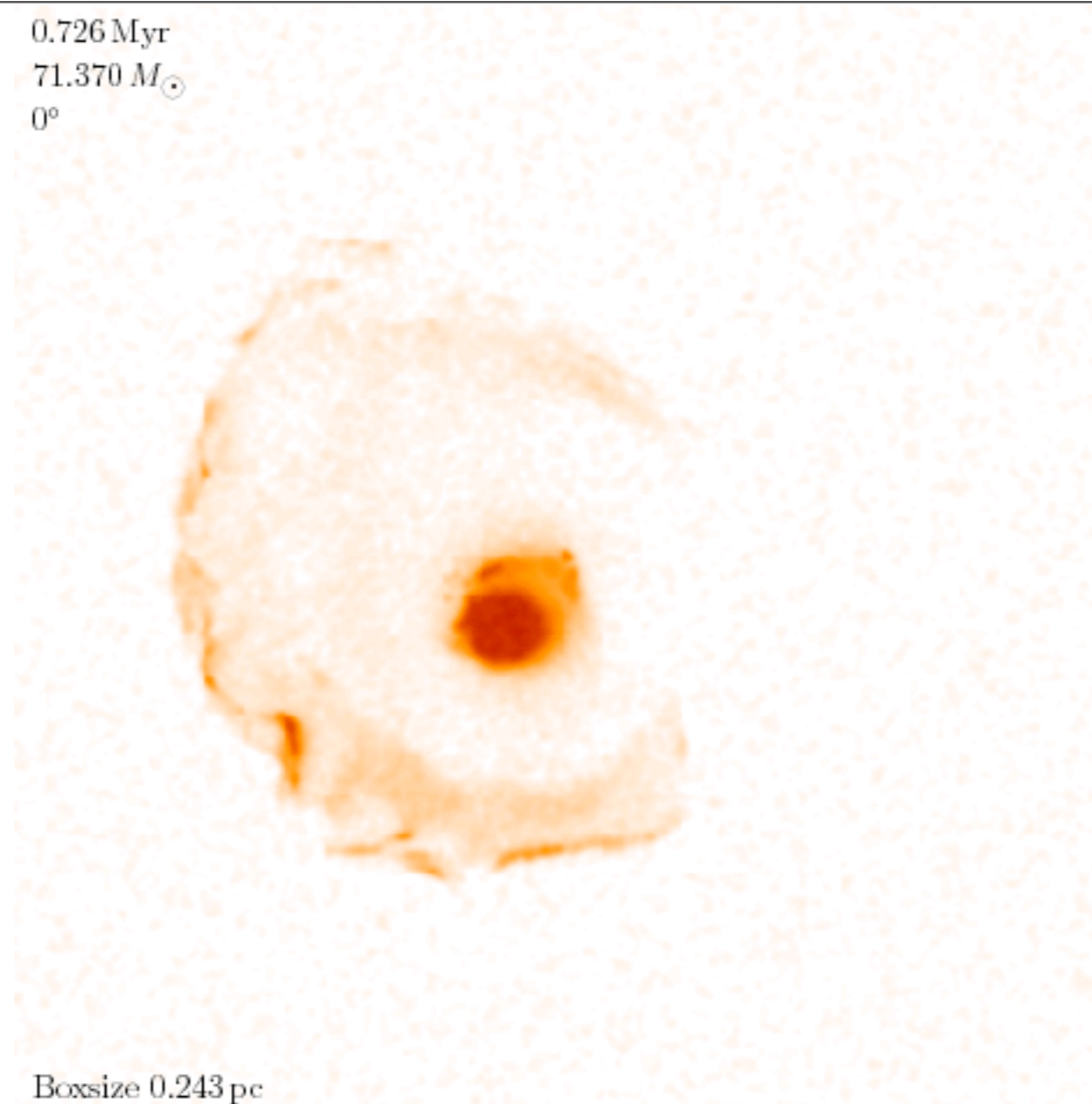


H II Region Morphologies



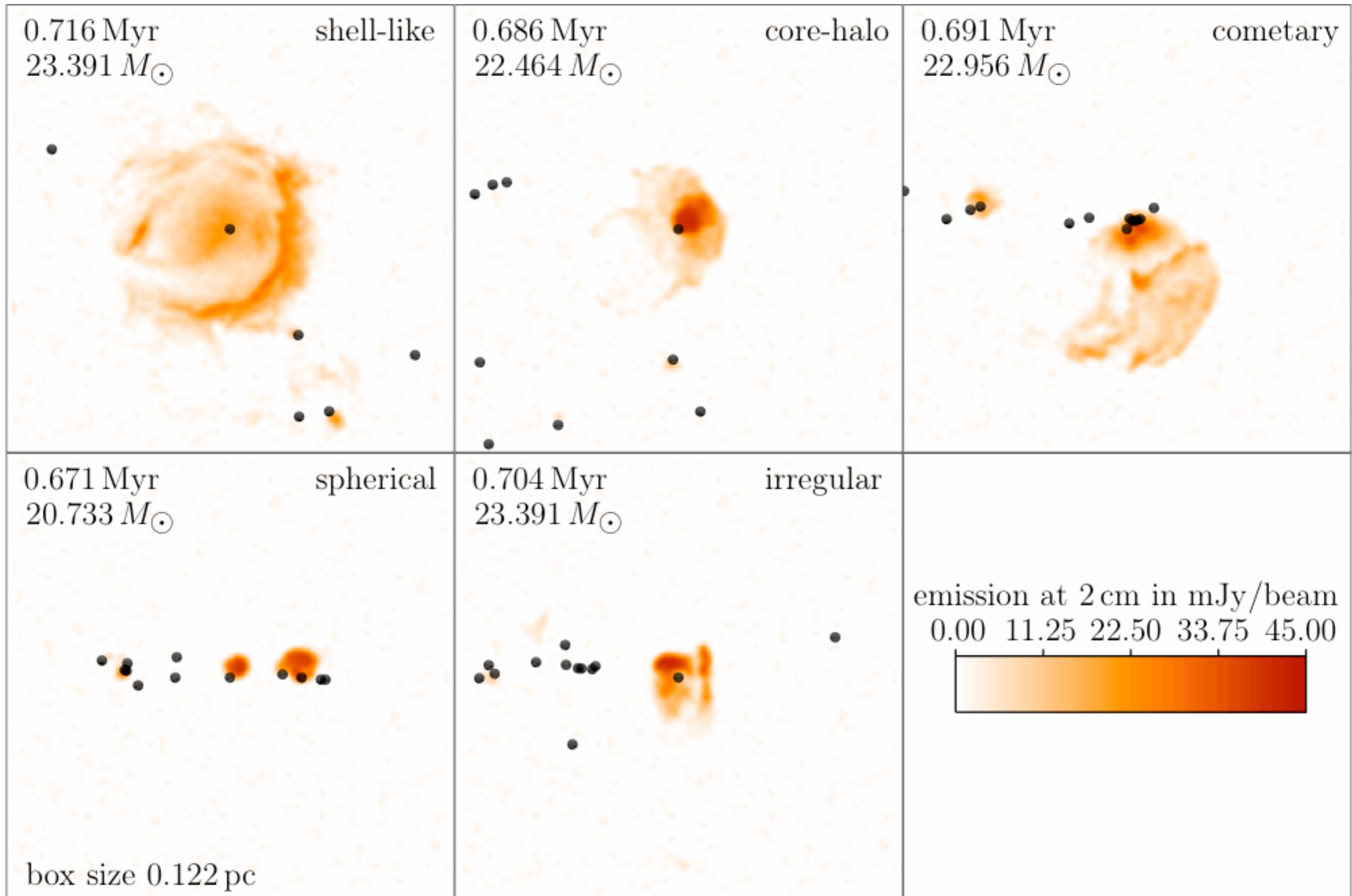
irregular morphology

H II Region Morphologies

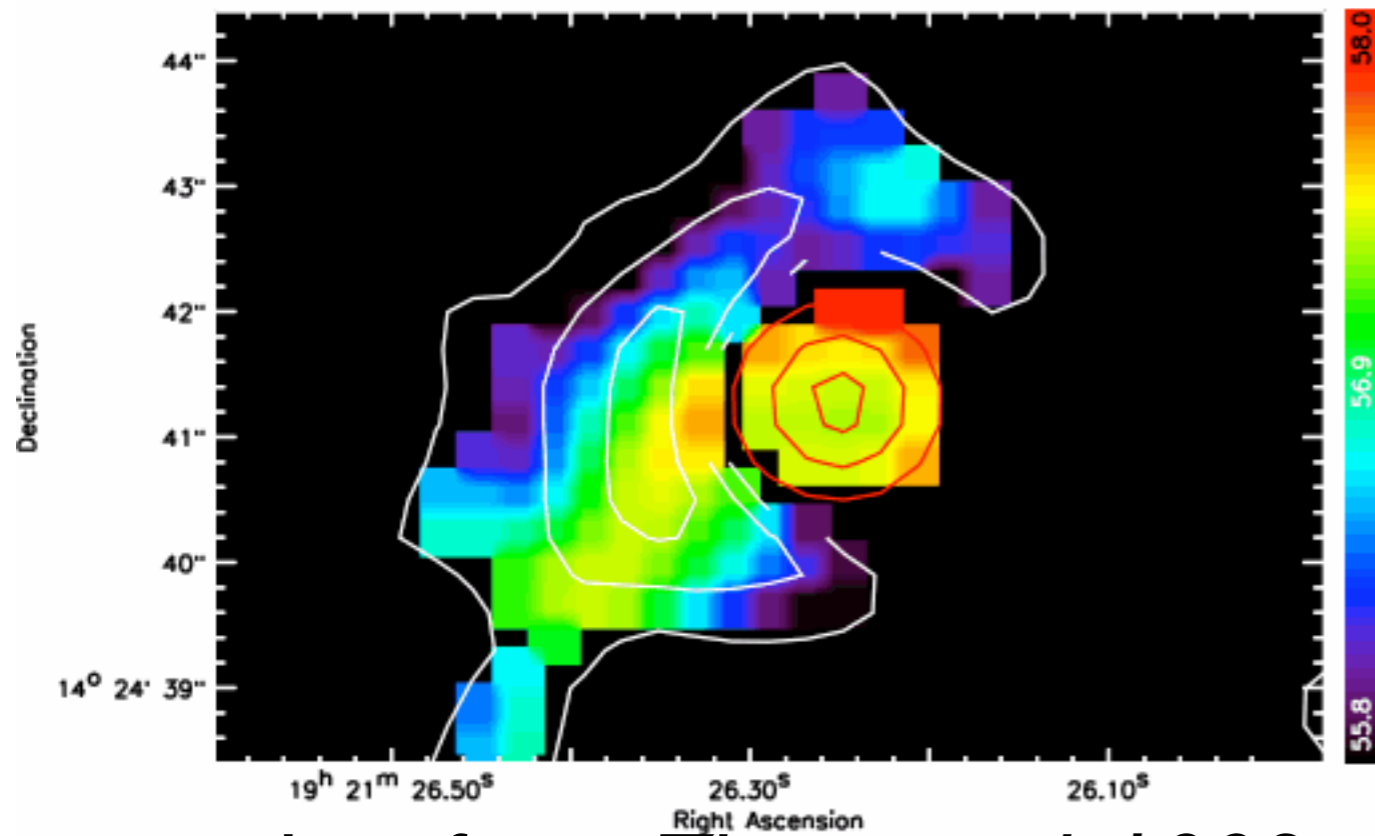
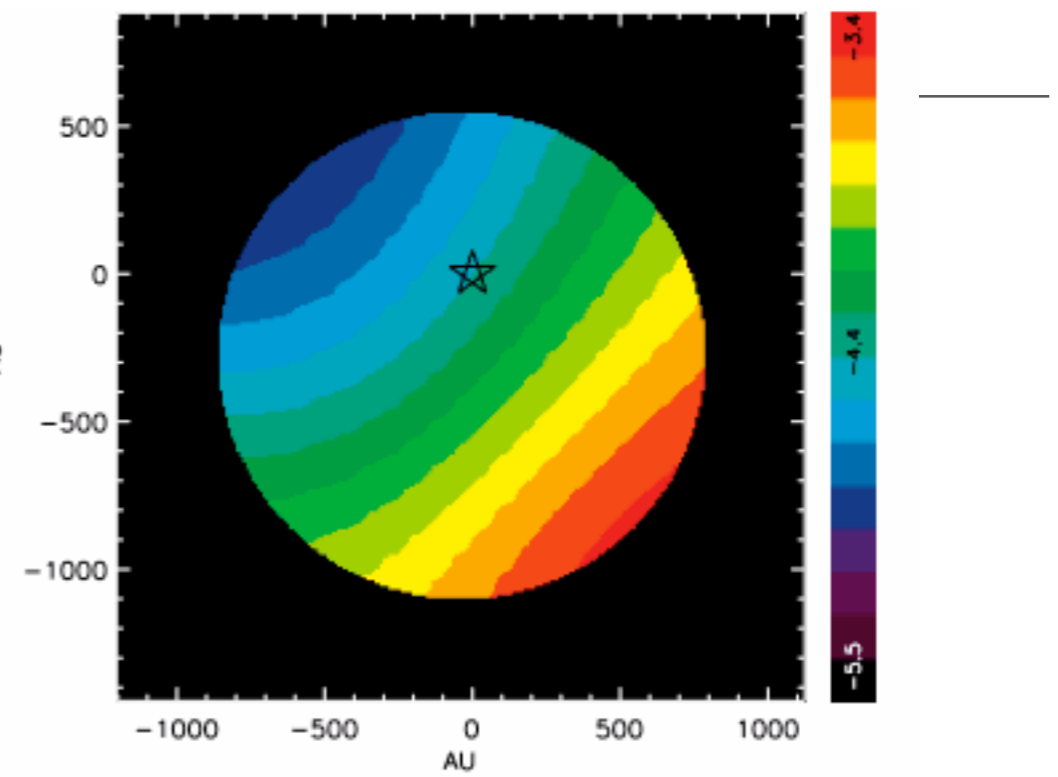
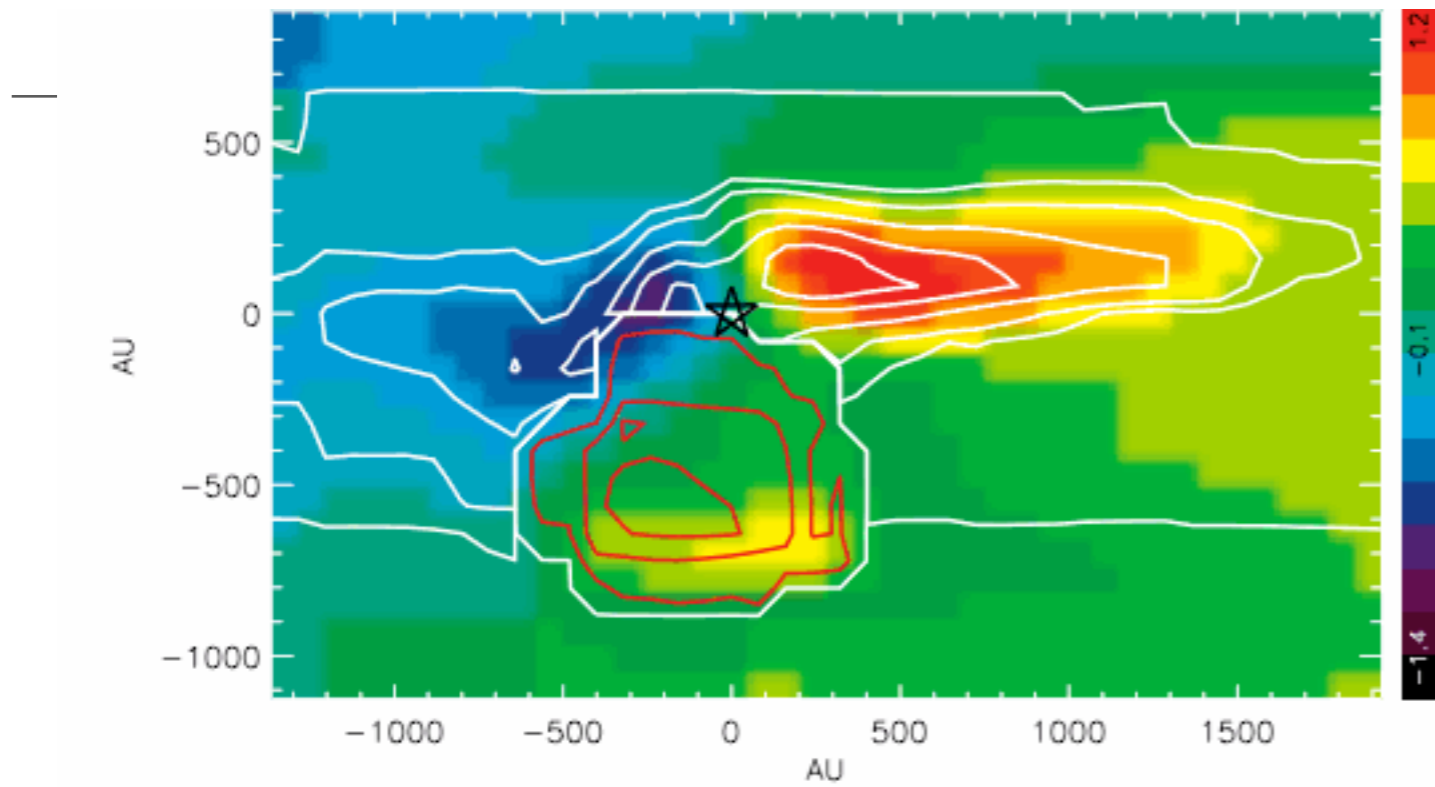


Morphology of HII region depends on viewing angle

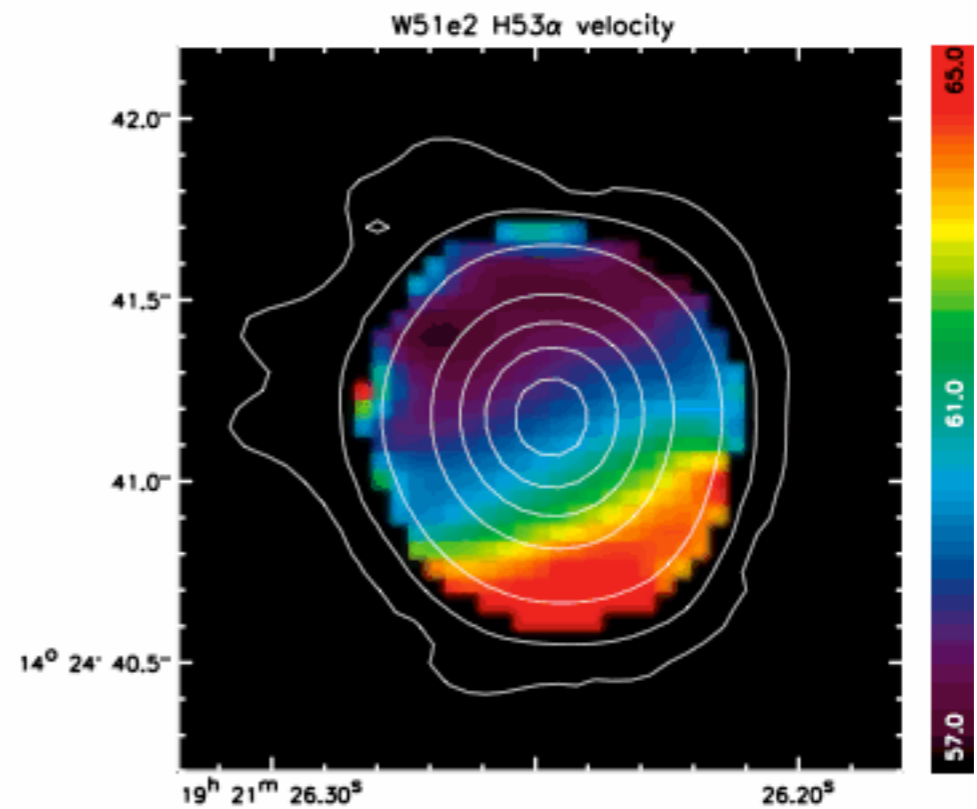
H II Region Morphologies



Comparison with observations: W51e2



data from *Zhang et al. 1998*



Keto & Klaassen 2008

Conclusions

- Molecular clouds can form at the cross section of **converging** flows by thermal instability
- MCs are **dynamic** objects with no distinct boundaries where warm and cold gas co-exist
- Ambipolar diffusion has only **little** influence on star formation
- Regions of **massive** star formation: rapid accretion through dense, unstable flows
- Ionization feedback does **not** shut off accretion
- HII regions are highly **variable** in time and shape
- All classified morphologies are seen in one run
- HII region will be gravitationally trapped (resolves lifetime problem, *Keto 2003*)